LOW TEMPERATURE SYNTHESIS OF THE MICROWAVE DIELECTRIC MATERIAL, BARIUM MAGNESIUM TANTALATE (BMT)

by

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Wireless communication systems utilize microwave dielectrics for coupling, selecting and filtering microwaves. Over the past several years there has been an increased demand for smaller, lighter and temperature stable devices. An important material that has been studied extensively for these applications is barium magnesium tantalate (BMT). Although BMT has very good dielectric properties: relatively high dielectric constant (25), temperature stability and low dielectric loss in the microwave region ($Q_{d*}f_o \approx 150,000$ GHz at 4.9 GHz), it can be expensive to produce because of the high sintering temperatures (>1600°C) required to obtain the desired properties.

The objective of this study was to dope BMT with ZnGa$_2$O$_4$, Ga$_2$O$_3$, and ZnO to try and lower the sintering temperature without severe degradation of the microwave dielectric properties. The study showed that the BMT co-doped with both Ga$_2$O$_3$ and ZnO gave the best properties at the lowest sintering temperatures. BMT co-doped with 4mol% Ga$_2$O$_3$ and ZnO has been successfully sintered at 1400°C for 2 hours had an average density of 95% with a dielectric constant of 24 and a $Q_{d*}f_o$ of 130,000 at 4.9 GHz. BMT co-doped with 8mol% Ga$_2$O$_3$ and ZnO and sintered at 1450°C for 2 hours had an average density of 94%, a dielectric constant of 24 and a $Q_{d*}f_o$ of 135,000 at 4.9GHz. The BMT materials doped with ZnGa$_2$O$_4$ and Ga$_2$O$_3$ both had average densities of over 95% and dielectric constants of approximately 24 but high dielectric
loss. The BMT doped with Ga$_2$O$_3$ had a Q$_{d^*f_0}$ of only 84,000 at 4.9 GHz and the BMT doped with ZnGa$_2$O$_4$ had a Q$_{d^*f_0}$ of 93,000 at 4.9 GHz. Phase evolution and densification behavior of these materials are described.

### DESCRIPTORS

| Barium Magnesium Tantalate | BMT |
| Cellular Phones | Dielectric Constant |
| Dielectric Resonator | Dopants |
| Liquid Phase Sintering | Microwaves |
| Ordering | Q-factor |
| Resonator | Solid State Sintering |
| Temperature Coefficient of Resonant Frequency | Zinc Gallate |
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# TABLE OF CONTENTS

<table>
<thead>
<tr>
<th>Section</th>
<th>Title</th>
<th>Page</th>
</tr>
</thead>
<tbody>
<tr>
<td>1.0</td>
<td>INTRODUCTION</td>
<td>1</td>
</tr>
<tr>
<td>2.0</td>
<td>BACKGROUND</td>
<td>3</td>
</tr>
<tr>
<td>2.1</td>
<td>Microwave Communications</td>
<td>3</td>
</tr>
<tr>
<td>2.1.1</td>
<td>Dielectric Resonators</td>
<td>4</td>
</tr>
<tr>
<td>2.1.2</td>
<td>Microwave Dielectric Properties</td>
<td>6</td>
</tr>
<tr>
<td>2.1.3</td>
<td>Microwave Dielectric Property Measurements</td>
<td>12</td>
</tr>
<tr>
<td>2.2</td>
<td>High Q Perovskite Ceramics</td>
<td>14</td>
</tr>
<tr>
<td>2.2.1</td>
<td>Barium Magnesium Tantalate (BMT)</td>
<td>15</td>
</tr>
<tr>
<td>3.0</td>
<td>OBJECTIVES</td>
<td>32</td>
</tr>
<tr>
<td>4.0</td>
<td>EXPERIMENTAL PROCEDURE</td>
<td>33</td>
</tr>
<tr>
<td>4.1</td>
<td>Synthesis of BMT Powders</td>
<td>33</td>
</tr>
<tr>
<td>4.2</td>
<td>Powder Processing</td>
<td>34</td>
</tr>
<tr>
<td>4.3</td>
<td>Microstructural Characterizations</td>
<td>35</td>
</tr>
<tr>
<td>4.3.1</td>
<td>X-ray Diffraction (XRD) Analysis</td>
<td>35</td>
</tr>
<tr>
<td>4.3.2</td>
<td>Scanning Electron Microscopy (SEM) Characterization</td>
<td>36</td>
</tr>
<tr>
<td>4.3.3</td>
<td>Density Measurements</td>
<td>36</td>
</tr>
<tr>
<td>4.4</td>
<td>Microwave Dielectric Property Measurements</td>
<td>37</td>
</tr>
<tr>
<td>4.4.1</td>
<td>Measurement of the Dielectric Constant</td>
<td>37</td>
</tr>
<tr>
<td>4.4.2</td>
<td>Measurement of the Quality Factor</td>
<td>38</td>
</tr>
</tbody>
</table>
5.0 RESULTS ................................................................................................................. 47

5.1 BMT Doped with 8mol% ZnO and Ga2O3 ............................................................... 47

5.1.1 Powder Development ......................................................................................... 47

5.1.2 Ceramics Sintered at a Constant Heating Rate .................................................. 47

5.1.3 Ceramics Sintered Isothermally at 1350oC .......................................................... 49

5.1.4 Ceramics Sintered at Various Temperatures for 2 hours .................................. 51

5.2 Comparing the Concentration of ZnO and Ga2O3 on the Sintering and Microwave Dielectric Properties of BMT ............................................................ 52

5.3 Comparing the Effect of Several Similar Dopants (ZnO and Ga2O3; Ga2O3; ZnO and ZnGa2O4) on BMT Sintered at 1450oC for 2 hours ........................................... 54

5.4 Overall Effect of Density on the Q-factor ............................................................... 56

5.5 Effect of Additives and Process Condition of Ordering .......................................... 57

5.6 The Effect of additives and Processing Condition on the Dielectric Constant ....... 57

6.0 DISCUSSION ........................................................................................................... 89

6.1 Phase Evolution .................................................................................................. 89

6.1.1 BMT Doped with ZnO & Ga2O3 and ZnGa2O4 .................................................. 89

6.1.2 BMT Doped with Ga2O3, ZnO and MgO ............................................................. 90

6.2 Densification Behavior ....................................................................................... 91

6.2.1 BMT Doped with ZnO & Ga2O3 and ZnGa2O4 .................................................. 91

6.2.2 BMT Doped with Ga2O3, ZnO and MgO ............................................................. 92

6.3 Q-factors ............................................................................................................ 94

6.3.1 BMT Doped with ZnO & Ga2O3 and ZnGa2O4 .................................................. 94

6.3.2 BMT Doped with each of the following Ga2O3, ZnO and MgO ......................... 94

7.0 CONCLUSION ..................................................................................................... 97
8.0 SUGGESTIONS FOR FUTURE WORK ................................................................. 98

Appendix A .............................................................................................................. 100
    Chemical Analysis of Starting Materials ............................................................ 100

Appendix B .............................................................................................................. 102
    Procedure for the Synthesis of BMT Ceramics .................................................. 102

Appendix C .............................................................................................................. 105
    Procedure for the Network Analyzer ................................................................. 105

BIBLIOGRAPHY ................................................................................................... 121
LIST OF TABLES

Table 1  Starting Materials for Various Compositions of BMT ......................... 46

Table 2  Ionic Radii of BMT and Other Relevant Ions ........................................ 96
LIST OF FIGURES

Figure 1  Microwaves are electromagnetic radiation. They lie in the region of $10^{-1}$ to $10^{-3}$ m ................................................................. 23

Figure 2  The microwave spectrum and some of the current applications utilizing microwaves. [64] .................................................................................. 24

Figure 3  Fields in a microwave dielectric resonator in the simplest standing wave mode (TE-O1d): (a) magnetic field, H; (b) electric field, E; (c) variation in Ef and Hz with r at Z=0 with reference to cylindrical coordinates ................................................................. 25

Figure 4  Mechanisms of dielectric polarization .......................................................................................................................... 26

Figure 5  A sample resonance curve that is produced when a DR is coupled to a microwave circuit .................................................................................................................................................. 27

Figure 6  The parameter used describes the signals that the network analyzer is measuring. The dielectric constant is measured by the $S_{21}$ parameter and the Q-factor is measured using the $S_{11}$ parameter ........................................................................................................................................ 28

Figure 7  The Q-factor of a microwave ceramic can be measured by using the single port reflection technique. The DR is placed on a piece of Styrofoam in the center of a silver cavity. Microwaves are sent from the network analyzer into the cavity. The signal bounces around the cavity and exits though the same port as it entered .................................................................................. 29

Figure 8  The Courtney method can be used to measure the dielectric constant of a microwave ceramic. The ceramic is placed between two metal plates. The signal comes from the network analyzer into the cavity and leaves the opposite side back into the network analyzer ........................................ 30

Figure 9  The perovskite structure of BMT. The barium ions occupy the corner positions of the cube. The magnesium and tantalum ions occupy the body-centered positions. And the oxygen anions occupy the face-centered positions .......................................................................................................................... 31

Figure 10  The procedure for the synthesis of BMT. When BMT was doped with ZnO, Ga2O3, and MgO, the dopant(s) was added before the first ballmilling. When the BMT was doped with ZnGa2O4, the dopant was added before the second ballmilling .................................................................................. 40

Figure 11  The JCPDS X-ray diffraction pattern for BMT (18-176) ........................................................................................................ 41
Figure 12  The JCPDS X-ray diffraction pattern for ZnGa$_2$O$_4$ (3-1155)………………………42

Figure 13  Courtney method set-up for the measurement of the dielectric constant……………43

Figure 14  Single port reflection technique set-up for the measurement of the quality factor. …………………………………………………………………………………………………44

Figure 15  A smith chart showing critical, over and under coupling…………………………45

Figure 16  X-ray diffraction pattern of calcined BMT powder doped with 8mol% ZnO and Ga$_2$O$_3$. This pattern shows only BMT peaks and no superlattice peaks……………………………58

Figure 17  SEM micrograph of calcined BMT powder doped with 8mol% ZnO and Ga$_2$O$_3$. There is no evidence of any second phases………………………………………………..59

Figure 18  The density Vs sintering temperature for BMT co-doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at a constant heating rate of 10$^\circ$C/min………………………………………..60

Figure 19  The Q*f vs sintering temperature for BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at a constant heating rate of 10$^\circ$C/min. The largest increase in Q-factor was between the temperatures of 1400$^\circ$C and 1450$^\circ$C; this corresponds to densities where the pores would be closing off. (The resonant frequency for these pellets was approximately 4.9 GHz.)…………………………………………………………………………61

Figure 20  Selected SEM micrographs of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at a constant heating rate of 10$^\circ$C/min. Some of the zinc and gallium rich second phase has been circled. The only pellet shown that did not show any evidence of Zn and Ga rich second phase is shown in figure (a) and was sintered at 1300$^\circ$C for 0 hours. The pellets were also sintered at 1400$^\circ$C for 0 hours (b), 1450$^\circ$C for 0 hours (c), 1500$^\circ$C for 0 hours (d), and1600$^\circ$C for 0 hours(e)…………………………………………………………………..62

Figure 21  An XRD pattern of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at a constant heating rate of 10$^\circ$C per minute increasing to 1500oC and then immediately decreasing at 10$^\circ$ per minute to room temperature………………………………………...63

Figure 22  The density vs. sintering time for BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1350$^\circ$C…………………………………………………………………………..64

Figure 23  The Q*f vs. sintering time for BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1350$^\circ$C (The resonant frequency for these pellets was approximately 4.9 GHz.)………65

Figure 24  Selected SEM micrographs of sintered surfaces of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered isothermally at 1350$^\circ$C. Some of the zinc and gallium rich second phase has been circled. The pellet shown in micrograph (a) has been sintered at 1350$^\circ$C for 0 hours and shows no evidence of any Zn and Ga rich second phase. The pellet shown in micrograph (b) has been sintered at 1350$^\circ$C for 2 hours, (c) has been sintered at 1350$^\circ$C for
6 hours. Micrograph (c) was taken using the backscattered electrons and therefore the Zn and Ga rich crystals appear darker than the BMT grains. And the pellet shown in micrograph (d) has been sintered at 1350°C for 24 hours.

Figure 25  A SEM micrograph of a fracture surface of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1350°C for 1 hour. Some of the Zn and Ga rich second phase has been circled.

Figure 26  A typical pattern from BMT doped with 8mol% ZnO and Ga$_2$O$_3$. This particular pattern came from a sample sintered at 1350°C/24hr. The circles peak has been identified as the start of barium tantalite. All other peaks have been identified as BMT peaks.

Figure 27  A graph of Density vs. Sintering Temperature (time=2 hrs) for BMT Doped with 8mol% ZnO and Ga$_2$O$_3$.

Figure 28  The Q$^*f$ vs. sintering temperature (time=2 hrs) for BMT doped with 8mol% ZnO and Ga$_2$O$_3$. (The resonant frequency was approximately 4.9 GHz.)

Figure 29  Selected SEM micrographs of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at various temperatures for 2 hours. Some of the zinc and gallium rich second phase has been circled. Micrograph a) has been sintered for 1450°C for 2 hours and is imaged by the backscattered electrons. Micrographs b) through f) have been sintered for 1600 °C for 2 hours.

Figure 30  An XRD pattern of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1500°C for 2 hours. The largest superlattice peak has been marked with an arrow.

Figure 31  An XRD pattern of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1600°C for 2 hours. This was taken on the top of the pellet; the bottom has the same peaks with the second phase peaks smaller. The zinc gallate peaks have been marked with an asterisk and the largest super lattice peak has been marked with an arrow.

Figure 32  A Comparison of the Q$^*f$ of 4mol% and 8mol% ZnO and Ga$_2$O$_3$ doped BMT. Although the BMT doped with 4mol% had slightly higher densities, the Q-factors were both higher in the BMT doped with 8mol% ZnO and Ga$_2$O$_3$.

Figure 33  Selected SEM micrographs of BMT doped with 4mol% (a) and b)) and 8mol% (c) and d)) ZnO and Ga$_2$O$_3$. Micrographs a) and c) have been sintered at 1400°C for 2 hours and micrographs b) and d) have been sintered at 1500°C for 2 hours. Some of the zinc and gallium rich second phase has been circled and the barium and tantalum rich second phase is indicated by a rectangle.

Figure 34  XRD pattern of the top of a sintered pellet of BMT doped with 4mol% ZnO and Ga$_2$O$_3$. The barium tantalate second phase peaks have been circled and the largest superlattice peak has been marked with an arrow.
Figure 35  An XRD pattern of the bottom of a pellet doped with 4mol% ZnO and Ga$_2$O$_3$ and sintered at 1500°C for 2 hours ................................................................. 77

Figure 36  An XRD of the top of a pellet of BMT doped with 4mol% ZnO and Ga$_2$O$_3$ and sintered at 1500°C for 2 hours with the outer surface removed ............................................ 78

Figure 37  An XRD of the bottom of a pellet of BMT doped with 4mol% ZnO and Ga$_2$O$_3$ and sintered at 1500°C for 2 hours with the outer surface removed ............................................ 79

Figure 38  An XRD pattern of a BMT pellet doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1400°C for 2 hours ................................................................. 80

Figure 39  A Comparison of the Q*f of BMT doped with combinations of ZnO and Ga$_2$O$_3$. ............................................................................................................. 81

Figure 40  SEM micrograph of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1450°C for 2 hours. Some of the zinc and gallium rich second phase has been circled .............. 82

Figure 41  SEM micrograph of BMT doped with 8mol% Ga$_2$O$_3$ and sintered at 1450°C for 2 hours. Some of the gallium rich second phase has been circled ........................................ 83

Figure 42  An XRD pattern of a BMT pellet (top) doped with 8mol% Ga$_2$O$_3$ and sintered at 1450°C for 2 hours. The largest superlattice peak has been marked with an arrow ............. 84

Figure 43  An XRD pattern of a BMT pellet (bottom) doped with 8mol% Ga$_2$O$_3$ and sintered at 1450°C for 2 hours. The largest superlattice peak has been marked with an arrow ............. 85

Figure 44  SEM micrograph of BMT doped with 8mol% ZnGa$_2$O$_4$ and sintered at 1450°C for 2 hours. Some of the second phase rich in gallium and zinc is circled and some of the second phase rich in barium and tantalum is indicated by a rectangle......................... 86

Figure 45  An XRD pattern of BMT doped with 8mol% ZnGa$_2$O$_4$ and sintered at 1450°C for 2 hours ............................................................................................................. 87

Figure 46  XRD pattern of the bottom of a pellet of BMT doped with 8mol% ZnGa$_2$O$_4$ and sintered at 1450°C for 2 hours. The zinc gallate second phase peaks have been marked with an asterisk and the largest superlattice peak has been marked with an arrow ............. 88

Figure 47  The S$_{21}$ spectrum ranging from 1 GHz to 10 GHz ................................................................. 116

Figure 48  The screen on the Courtney program where the actual data is entered so the program can calculate the dielectric constant ................................................................. 117

Figure 49  An example of the results obtained from the Courtney Software ........................................ 118
Figure 50  The resonant coupling circle on the Smith chart for measuring the Q-factor using the single-port reflection technique .................................................................119

Figure 51  An example of the results obtained from the single-port reflection technique for measuring the Q-factor .................................................................................................120
1.0 INTRODUCTION

Recently wireless communication systems using microwaves to carry information have become increasingly popular. Microwaves are able to carry more information than radio waves because of their higher frequency. These systems include cellular phones and satellite communications. Wireless communications utilize microwave dielectric components such as dielectric resonators for coupling, selecting and filtering microwaves. Traditional wiring and transistors will not work in these systems due to lead reactance and transit time.

There are three main, very important properties of a dielectric resonator material: relatively high dielectric constant, low dielectric loss, and temperature stability. A relatively high dielectric constant enables the dielectric resonator to be smaller since the dielectric constant is inversely proportional to the size. Low dielectric loss or high quality factor gives a smaller bandwidth at the resonant frequency, a lower degree of noise, and less power loss. Temperature stability is important so that the device can be used in any environment without detrimental effects on the signal.

There are many materials currently used in these applications, many of which are ceramics with a perovskite structure. The two most commonly utilized materials are barium zinc tantalate (BZT) and barium magnesium tantalate (BMT). Although these materials have very good microwave dielectric properties, they can be very expensive to produce. BMT and BZT require high sintering temperatures, often greater than 1600°C, and long annealing times, sometimes for 24 hours or more to obtain the desired properties. This study focused on lowering the sintering temperature of BMT, through the addition of dopants, while maintaining the dielectric properties of the material. BMT was doped with ZnGa$_2$O$_4$ (8mol%); Ga$_2$O$_3$ (8mol%);
ZnO (8mol%); and Ga$_2$O$_3$ along with ZnO (4mol% and 8mol%). The ceramics were synthesized through a solid state reaction and sintered at various temperatures between 1300°C and 1600°C for various times. The dielectric properties of the materials were measured using a network analyzer. The dielectric constant was measured by the Courtney method and the loss was measured by the single port reflection technique. The microstructure analysis was done by x-ray analysis and scanning electron microscopy.

The major technical objective of this research was to investigate the synthesis of BMT using the conventional mixed oxide route, and identify sintering additives that could lower the sintering temperature while maintaining the dielectric properties.
2.0 BACKGROUND

2.1 Microwave Communications

Electromagnetic radiation is a wave of oscillating electric and magnetic influences. The electromagnetic spectrum ranges from 1 MHz to $10^6$ GHz and can be broken down into ranges as shown in figure 1. The radio frequency is in the range of 300 KHz to 300 MHz. Radio waves are used to carry information in AM and FM broadcast radio, short wave radio, and VHF television channels. These signals can be carried using traditional transistors, tubes and circuits. Microwaves can also be used to carry information. Microwaves lie in the region from 300 MHz to 300 GHz. The microwave region can be broken down further into the ultra high frequency (UHF) region from 300 MHz to 3 GHz, the super high frequency (SHF) region from 3 GHz to 30 GHz, and the extremely high frequency (EHF) region from 30 GHz to 300 GHz. Microwaves are sometimes more useful in carrying information than radio waves because they have a higher frequency and shorter wavelength. The higher frequency gives a wider bandwidth capacity, enabling it to carry more information. The short wavelength allows the energy to be concentrated into a small area; this is the reason that microwave ovens are so convenient. There are many current applications that use microwaves, including cellular phones (900 MHz) and wireless internet (2.4 GHz). Figure 2 shows the microwave spectrum and many of the current applications. Most microwave systems are located below 30 GHz (between 300 MHz and 30 GHz) because at many frequencies above 30 GHz there is high atmospheric absorption making
them impractical for long-range communication and radar systems. There are some windows above 30 GHz that can be utilized; missile seekers are used at a frequency of about 90 GHz.

Due to lead reactance and transit time microwave systems require special devices to carry and filter the electromagnetic waves; conventional transistors, IC’s and wiring will not work. The inductive reactance of the wire increases with frequency and is negligible when the frequencies are low. However, at microwave frequencies, the lead reactance is so high that almost all of the oscillator voltage is dropped across the connecting wire and never gets to the load resistor, meaning that most of the signal (information) is lost.

2.1.1 Dielectric Resonators

A resonator is the most commonly used configuration for filters and oscillators. At high frequencies (>3GHz) the resonator is made from a piece of dielectric rod, called a dielectric resonator. In 1939, Richtmyer\(^{(1)}\) showed that unmetallized dielectric objects could function similarly to metallic cavity resonators; he called these dielectric resonators (DR’s). A DR, usually a ceramic, is able to resonate at the frequency of the carrier signal to allow that signal to be efficiently separated from other signals in the microwave band; this frequency is called the resonant frequency \(f_o\). The resonant frequency depends on the dielectric material and the size of the resonator.

It was not until the late 1960’s that dielectric resonator ceramics were developed for resonator applications, before this, metallic high Q cavities were utilized. During this time there were many pioneers investigating the behavior of dielectrics at microwave frequencies, including: Rupprecht\(^{(2)}\), Spitzer\(^{(3)}\), Hakkl\(^{(4)}\), Coleman\(^{(4)}\), and Cohn\(^{(5)}\). At this time the
applications were limited due to a lack of suitable materials. Rutile phase (TiO$_2$) was often used but its large resonant frequency fluctuations with temperature made it impractical for many applications. Single crystal alumina (sapphire) was also used but it is very expensive. By the 1970’s, research and development for temperature stable dielectric resonator materials had begun worldwide.

There are many advantages of ceramic dielectric resonators compared to traditional resonators including: size reduction (higher dielectric constant), greater degree of circuit and subsystem integration (due to simpler coupling schemes), better circuit performance (with regard to temperature changes and losses), and a reduction in the overall cost of the systems. Over the years the need for smaller, lighter, faster and cheaper communication devices has driven the industry to explore many new materials and systems. Today microwave dielectric resonators are used to create, filter, and select frequencies in oscillators, amplifiers and tuners.

The most common/advantageous dielectric resonator is a solid, cylinder shape that can sustain an electromagnetic standing wave within its volume because of the reflection at the dielectric-air interface. Several DR’s can be assembled into devices to filter microwaves. These devices pass a range of frequencies and reject all others. Since the filter passes only certain frequencies there needs to be large out of band attenuation, low in band insertion loss, and a small bandpass to give greater selectivity. There are two mechanisms of selecting the frequency: absorption and reflection. In absorption, the frequency is selected by absorbing the out of band frequencies. In reflection, the frequency is selected by reflecting out the band frequencies and allowing only in band frequencies. There are three kinds of resonant modes that are often used for dielectric resonant filters: TEM mode, TE$_{011}$ and TM$_{010}$ or HE mode. The resonant modes describe the electric field lines around the resonator. The electric field lines of each mode are
shown in figure 3. The TE (transverse electric) has no transverse component of the electric field; the TM (transverse magnetic) has no transverse component of the magnetic field; and the HE (hybrid electro-magnetic) has a transverse component in the electric and magnetic fields.

There are three important properties that need to be considered for materials used for dielectric resonators. The temperature coefficient of resonant frequency, $\delta f$, describes the temperature stability of the material. This should be as close to zero as possible. The unloaded quality factor, $Q_0$, is inversely proportional to the losses of the material; therefore this should be as high as possible to get the minimum loss. And the dielectric constant, $K'$, ultimately determines the resonator dimensions and should also be as high as possible. These qualifications will be discussed separately and in more detail in the following sections.

2.1.2 Microwave Dielectric Properties

2.1.2.1 Dielectric Resonant Frequency. The resonant frequency is the frequency at which the dielectric resonator is utilized. The resonant frequency of a dielectric resonator in the TE01δ mode can be estimated by the following equation:

$$f_o \,(\text{GHz}) \approx \frac{8.553}{((K')^{-2}((\pi/4)(D^2L^2)))^{1/3}}$$

(2-1)

$f_o$ is the resonant frequency, $K'$ is the dielectric constant, $D$ is the resonator diameter in inches, and $L$ is the resonator height in inches. This equation gives the resonate frequency within approximately 6%.$^{(6)}$

The resonate frequency can be changed by altering the dimensions of the DR; from a manufacturing standpoint the height is the easiest dimension to change. It can also be changed
by perturbing the fringing fields outside the dielectric resonator with a metallic or dielectric tuner. A metal tuning stub will pull the frequency up and dielectric-tuning stub will push the frequency down. A draw back to using a tuning stub is that it will slightly decrease the Q-factor and slightly disturb the temperature stability.

2.1.2.2 **Dielectric Constant.** When a parallel-plate capacitor charged with a voltage has a dielectric between the plates, the dipoles of the material will align with the field. The fields of the dipoles oppose the capacitor field. The result is a reduction in the net electric field between the plates. The dielectric constant is defined as the ratio of capacitance of a capacitor with a dielectric between the plates to that with a vacuum between the plates. This is defined in equation 2-2.

\[ K' = \frac{C}{C_o} = \frac{\hat{\varepsilon}A/d}{(\hat{\varepsilon}_o A/d)} = \frac{\hat{\varepsilon}}{\hat{\varepsilon}_o} \]  

(2-2)

\( K' \) is the relative dielectric constant, \( C \) is the capacitance of the capacitor with the dielectric between the plates, \( C_o \) is the capacitance of the capacitor with a vacuum between the plates, \( A \) is the area of the plate of the capacitor, \( d \) is the distance between the plates, \( \hat{\varepsilon} \) is the permittivity of the material, and \( \hat{\varepsilon}_o \) is the permittivity of free space. The charged species of the material are displaced in response to the applied field; this increases the ability of the plates to store charge. This orientation of the charged species “ties up” charges on the plates of the condenser and thus neutralizes part of the applied field. The charge that isn’t neutralized (free charge) produces an electric field and voltage towards the outside. Therefore, a smaller external field is required to maintain the same surface charge because some of the charge is held by the polarization in the dielectric. There are four primary mechanisms of polarization: electronic (due to the shift of the electron cloud), atomic (due to the displacement of the positive and negative ions with respect to each other), dipole (due to the molecules containing permanent dipole rotating or rotation of
dipoles between two equivalent equilibrium positions) and interfacial (due to mobile charge carriers that are impeded by a physical barrier). The four primary mechanisms for polarization are shown in figure 4. Each of these involves a short-range motion of charge and contributes to the total polarization of the material.

The dielectric constant of the dielectric resonator material is important because it will ultimately determine the dimensions of the resonator. A cylindrically shaped DR sustains an electromagnetic standing wave within its volume because of the reflection at the dielectric air interface. The wavelength of the standing wave is about equal to the diameter of the cylinder. This is shown in equation 2-3 where D is the diameter of the resonator, C is the capacitance of the resonator, \( f_o \) is the resonant frequency and \( \hat{\epsilon}_r \) is the permittivity of the resonator.

\[
D = \frac{C}{f_o \hat{\epsilon}_r^{1/2}}
\]  

(2-3)

It is therefore important to have the dielectric constant as high as possible in order to reduce the size of the DR.

2.1.2.3 Quality factor. The quality factor or Q-factor is the ability of a dielectric resonator to store microwave energy with minimal signal loss. A perfect dielectric has no losses since its electrical conductivity is zero but real dielectric materials are partially conducting. Therefore the electromagnetic signal loses power as it goes through a dielectric material. The Q-factor is approximately equal to the inverse of the loss tangent; therefore for DR applications it is important for the losses to be minimal and the Q-factor to be as high as possible.

The Q-factor is a measure of the energy loss per cycle as compared to the energy stored in the fields outside the ceramic. The Q-factor is a dimensionless number that can be expressed as a ratio shown in equation 2-4.

\[
Q = \frac{2\pi \text{ (max. energy stored per cycle)/}}{\text{(average energy dissipated per cycle)}}
\]  

(2-4)
This ratio must be evaluated at the resonant frequency of the material and if the material has more than one resonant frequency, each may have a different Q-factor. When a DR is coupled to a microwave circuit a resonance curve can be produced (a sample curve is shown in figure 5) and by using this curve along with equation 2-5 the Q-factor can be calculated;

\[ Q = \frac{f_o}{\Delta f} \]  

\( f_o \) is the resonant frequency and \( \Delta f \) is the half power bandwidth at the resonant frequency. A high Q-factor material has a narrow bandwidth at the resonant frequency. This provides more frequency selections in a frequency range and a better signal to noise ratio. (A more detailed description of measuring Q-factors is given in section 4.4.2 and Appendix C.) This Q-factor is called the unloaded quality factor, \( Q_u \). For cavity resonators there are a number of factors that affect the \( Q_u \) including: dielectric loss of the dielectric itself, conduction in the walls of the enclosure, and radiation. These are usually referred to as \( Q_d \), \( Q_c \) and \( Q_r \) respectively. The losses of the dielectric itself can be broken down further into intrinsic and extrinsic losses. The intrinsic losses are dependant on the composition of the material. The extrinsic losses are determined by the microstructure of the material. The losses arise from interactions of phonons, in ideal crystals, and impurities, vacancies, grain boundaries, porosity and dislocations in real materials. The intrinsic losses are due to decay of the microwave into optical phonons and finally into thermal phonons.\(^{(7)}\) From theory it would seem that the highest Q-factor would be obtained from a single or perfect crystal, since in general, defects and impurities will increase phonon scattering, but there has not been any evidence from research to prove this. The Q-factor is often reported as the Q-factor times the resonant frequency (\( Q^*f \)). Theoretically the Q-factor decreases linearly with increasing frequency; comparing just the Q-factors of different resonators would not be a fair comparison. It is therefore customary to report Q-factors as \( Q^*f \). If it isn’t
clear which value is being reported, the Q-factor is dimensionless and the \( Q^* f \) value will have the units of the resonant frequency, usually of GHz.

There are three main mechanisms for microwave absorption: 1) losses in perfect (ideal) crystals because of anharmonic lattice forces which mediate interaction between the crystal and phonons, 2) losses in real but homogenous crystals caused by periodicity defects (isotopes, dopant atoms, vacancies, etc.) and 3) losses in real, inhomogeneous ceramics by extended dislocations, grain boundaries, vacancies, inclusions and second phases. The losses from point 3) are small and can usually be neglected at high frequencies.\(^{(7)}\) Points 1) and 2) can be seen as energy transfers from the excited microwave to transverse optical phonons. These optical phonons can then generate thermal phonons through scattering by impurities and through interaction with other phonons. This leads to dampening of the microwave. The dielectric loss for these scattering effects can be calculated by the following equation:

\[
\tan \delta = \omega \Sigma \left( \Delta \varepsilon / \varepsilon(0) \right) \left( \gamma \omega^2 \right)
\]

where \( \omega \) is the frequency and \( \gamma \) is the dampening constant. A linear increase of loss with frequency is therefore characteristic for these phonon effects. Therefore, the \( Q^* f_o \) is constant in the absence of extrinsic losses. And when there are extrinsic losses the \( Q^* f_o \) may not increase linearly with frequency.

2.1.2.4 Temperature Coefficient of Resonant Frequency. The temperature coefficient of resonant frequency, \( \tau_f \), determines how well a resonator will function when there are fluctuations in temperature. This includes the ambient temperature as well as heating by microwave absorption in the resonator structure. The \( \tau_f \) is a combination of three independent factors: temperature coefficient of the dielectric constant (\( \tau_\varepsilon \)), thermal expansion of the material (\( \alpha_L \)) and
the thermal expansion of the environment in which the resonator is mounted. This is shown in equation 2-7(6).

\[ \tau_f = \frac{1}{2} \tau e - \alpha_L \]  (2-7)

Changes in temperature can cause changes in the dielectric constant, the size of the resonator, and the resonant frequency. In order for materials to be used as oscillators and filters, it is important that the resonant frequency not shift significantly between -20°C and +80°C. If the resonant frequency is too dependent on the temperature, the carrier signal will drift in and out of resonance on hot and cold days. If these fluctuations cannot be avoided, it is preferred that the resonant frequency changes are small and occur linearly because then there are ways to compensate. A material with a small positive variation can often be combined with a material with a small negative variation to make the resonator stable. The equation for calculation is shown in equation 2-8 where \( \tau_k \) is the temperature coefficient of the dielectric constant and \( \alpha \) is the thermal expansion coefficient of the dielectric.

\[ \tau_f = \frac{-1}{2} \tau_k - \alpha \]  (2-8)

For DR applications the \( \tau_f \) should typically be equal to or less than ±3ppm/°C. The \( \tau_f \) can also be expressed as a function of frequency as shown in the following equation,

\[ \tau_f = \frac{\Delta f (\text{Hz})}{f_o (\text{MHz})}/\Delta T \]  (2-9)

where \( \Delta f \) is the resonant frequency at 60°C minus the resonant frequency at 25°C, \( f_o \) is the resonant frequency at 25°C and \( \Delta T \) is 35°C (60-25).
2.1.3 Microwave Dielectric Property Measurements

2.1.3.1 Network Analyzer. A network analyzer is an instrument that is able to measure the reflection and transmission characteristics of devices by applying a known swept signal and measuring the response of the device. The signal that is transmitted or reflected from the device is compared to the incident signal generated by a swept RF source. The signals are applied to a receiver for measurement, signal processing and display.

The network analyzer is able to measure many parameters. Figure 6 shows some of the possible parameters. These describe the signals that the network analyzer is measuring. For example, for a transmission method measurement ($S_{21}$), the electric field of the microwave signal entering the component input would be compared to the electric field of the microwave signal leaving the component output. Or more simply, the signal entering one side of a cavity would be compared to the signal collected on the other side of the cavity.

2.1.3.2 Measurement of the Quality Factor. There are many ways to measure the Q-factor of dielectric ceramics. A common method and the method used in this study, is the single-port reflection technique. The single port reflection technique was developed by Engineers at TransTech, Inc.\(^8\) This method uses the TE-01å mode and the $S_{11}$ parameter. The $S_{11}$ parameter is when the electric field of the microwave signal leaving the component input is divided by the electric field of the microwave signal entering the component input. (The signal is entering and leaving the cavity through the same port.) This technique uses a silver plated cavity that is at least three times the size of the DR. The DR must be separated from the walls of the cavity to reduce conduction loss; therefore the DR is placed on a piece of Styrofoam or other low loss
material. This set up is shown in figure 7. This technique uses the TE-01ä mode, the lowest and most easily coupled mode.

2.1.3.4 **Measurement of the Dielectric Constant.** The Courtney method is the most common method for measuring the dielectric constant of dielectric ceramics. In 1960, Hakki and Coleman\(^{(4)}\) introduced the parallel-plate dielectric resonator to measure the dielectric constant of ceramics. Courtney\(^{(9)}\) then analyzed the error and investigated the temperature effects of this technique and it is now known as the Courtney method. This is the method that was used to measure the dielectric constant of the materials in this study. The Courtney method is a transmission type measurement that uses a small brass cavity. The sample is placed in the center of the cavity away from the sides. The set up for the Courtney method is shown in figure 8. This technique uses the S\(_{21}\) parameter (the electric field of the signal entering the component input is divided by the electric field of the signal leaving the component output) and the TE-011 mode. The resonant frequency at the TE-011 mode, 3dB half power bandwidth, and the insertion loss are obtained from the network analyzer. This data along with the dimensions of the DR are entered into the Courtney software (developed by GDK products, Cazenovia, NY), which calculates the dielectric constant.

2.1.3.5 **Measurement the Temperature Coefficient of Resonant Frequency.** A simple technique used to measure the temperature coefficient of resonant frequency is to note the shift in frequency with respect to temperature. The DR is placed in the cavity and weakly coupled so that the tuning effects due to expansion are minimized. The set up is heated slowly using a hotplate and the resonant frequency is noted at 25\(^{\circ}\)C and 60\(^{\circ}\)C. Equation 2-8 can then be used to calculate the \(\delta_f\).
2.2 High Q Perovskite Ceramics

High Q Perovskite ceramics have the general formal $A(B'_{1/3}B''_{2/3})O_3$ where $A$ is a cation with a large ionic radius ($Ba^{+2}$, $Sr^{+2}$, ect.), $B'$ is a divalent cation with a smaller ionic radius ($Mg^{+2}$, $Zn^{+2}$, ect.) and $B''$ is a pentavalent cation with a smaller ionic radius ($Ta^{+5}$, $Nb^{+5}$, ect.). There are many different compounds with this structure that are utilized or have been investigated for DR applications; the most common are barium magnesium tantalate (BMT) and barium zinc tantalate (BZT). Table 1 shows a list of some common microwave dielectric resonator materials. Recently there has been an increased interest in the niobium compounds especially barium zinc niobate (BZN) because of a price increase in the already expensive tantalum oxide.

Since there are two different B-site cations, there are two different structures that can exist, the disordered state and the ordered state. The disordered state is a cubic structure; this is the same structure as $BaTiO_3$. The ordered state is a hexagonal perovskite, shown in figure 9. Galasso$^{(10-13)}$ investigated the structure of many perovskites with a 1:2 mixture of divalent and pentavalent cations in the B-sites. These materials can adopt a structure in which the two cations order onto individual (111) crystallographic planes in a two to one ratio. The extent of this ordering depends on the differences in the ionic radii and the charges of the B-site cations; the greater the differences the greater the ordering. In general, for charge differences $\geq 4$ the cations will order, but when the difference is less than 4 it will be disordered unless there is a large difference in the radii.
2.2.1 Barium Magnesium Tantalate (BMT)

As mentioned earlier, BMT is a perovskite ceramic that has been investigated extensively\(^\text{(14) (15) (16) (17) (18) (19) (20) (21) (22) (23) (24) (25) (26) (27) (28) (29) (30) (31) (32) (33) (34) (35) (36) (37) (38) (39) (40) (41) (42) (43) (44) (45) (46) (47) (48) (49) (50) (51) (52) (53) (54) (55) (56) (57)}\) and is utilized for DR applications. BMT has a reported dielectric constant of approximately 25 and a temperature coefficient of resonant frequency of approximately zero; these values have been reported consistently among research groups. There is much less agreement concerning the dielectric loss, reported Q*\(f\) values range from 10,400 GHz\(^\text{(58)}\) to 600,000 GHz even though this is supposed to be a material constant. Some of these variations may be due to different measurement techniques but most are probably due to variations in extrinsic losses due to defects introduced in processing and the purity of the starting materials. Although several factors have been suggested, including B-site cation ordering and point defects, the ranking of the important sources of losses are not clear.

Although BMT has excellent dielectric properties, it can be very expensive to produce. BMT is very refractory and therefore requires very high sintering temperatures (1550\(^\circ\)C to 1650\(^\circ\)C) and often long annealing times to get the desired properties. Many research groups have investigated solutions to these problems including: chemical synthesis and sintering aids. Unfortunately, these attempts tend to significantly degrade the dielectric properties.

2.2.1.1 Synthesis of BMT. The most common synthesis method for BMT and the method used in this study is through a solid-state reaction. The metal oxides are used as starting materials and mixed homogeneously by ballmilling. The powders are then heated to a high temperature to form BMT according to the following reaction:

\[
3\text{BaCO}_3(s) + \text{Ta}_2\text{O}_5(s) + \text{MgO}(s) \xrightarrow{1350^\circ\text{C/5 hr}} \text{Ba}_3\text{Ta}_2\text{MgO}_9(s) + 3\text{CO}_2(g)
\]
This is currently the most common synthesis method because it tends to give ceramics with the best dielectric properties, although they require very high sintering temperatures (1600°C or greater) and sometimes long annealing times (32 hours or longer). (18, 24, 29, 59, 60)

Another synthesis method is sol-gel synthesis. Chemical synthesis seemed attractive because it tends to produce powders that are more homogeneous with a finer particle size and better control over the particle morphology that often gives improved sinterability. Several research groups were able to produce ceramics with relative densities, over 96%, at temperatures between 1300°C and 1400°C for as little as two hours. Renoult et. al. (43) was even able to get complete B-site ordering at 1400°C for 5 hours. (As mentioned earlier, some consider B-site cation ordering to be very important for achieving low loss ceramics.) Despite the high density and ordering parameter the Q-factor was only 6,410 at 7.71GHz (Q*\(f\) = 49,421 GHZ). Even when the sintering temperature of these ceramics was raised to 1600°C, the losses are still relatively high. This may be due to the higher point defect concentrations inherent to low temperature, chemically prepared materials or undetectable second phases.

Maclaren et. al. (35) synthesized BMT through hydrothermal synthesis. In hydrothermal synthesis, ceramic sols are created by chemical reactions in an aqueous or organo-aqueous solution under the simultaneous application of heat and pressure. This is usually done in the presence of an alkali or acid that acts as a catalyst. There are some advantages over sol-gel synthesis. Hydrothermal synthesis uses simple salts such as acetates and nitrates instead of expensive alkoxides, the materials are produced directly so there is no need for a calcination step and the particles are produced in a sol form so these sols can be used directly in the production of green bodies without requiring the conversion to a dry powder. This method produced BMT at 200°C. The dielectric properties were not measured but there was evidence of barium carbonate
second phases and particles that were deficient in magnesium, both of which suggest that the
dielectric properties would not be acceptable.

Lu et. al\(^{(37)}\) produced BMT through homogeneous precipitation. This involves the
simultaneous precipitation of the ions to give better mixing. One of the problems that occurred
was contamination of chlorine ions (from the TaCl\(_5\)) that can cause cracking, and the sintering
temperature was still 1690°C. None of the dielectric properties were reported.

Another method commonly investigated to reduce the sintering temperature of BMT is
the addition of sintering aids. Nomura et. al\(^{(39)}\) added 1mol% Mn and got a Q of 16,800 at 10.5
GHz (Q\(^*\)f = 176,400) but the ceramics were sintered at 1550°C for 3 hours. Kim et. al\(^{(31)}\) tried
0.083 mol% Ca and got a Q of 9,000 at 10.5 GHz (Q\(^*\)f = 94,500 GHz) by sintering at 1650°C for
2 hours and then annealing at 1500°C for 2 hours. Katayama et. al\(^{(27)}\) added Co to make
Ba(Co\(_{0.125}\)Mg\(_{0.875}\)\(_{2/3}\)Ta\(_{2/3}\)O\(_3\), he reported a Q\(^*\)f of 103,100GHz but sintered the ceramics at
1600°C for 64 hours. Furuya et. al\(^{(23)}\) doped BMT with W to make 0.5mol% Ba(Mg\(_{1/2}\)W\(_{1/2}\)O\(_3\). He reported a Q of 40,000 at 10 GHz (Q\(^*\)f = 400,000) but he sintered the ceramics between
1500°C and 1600°C for 24 hours. Chen et. al\(^{(21, 22)}\) tried doping with 3.0 wt% NaF and 5%
BaO:Nd\(_2\)O\(_3\):5TiO\(_2\), these samples were sintered between 1200°C and 1400°C but were reported to
have low Q-factors. As mentioned earlier, many researchers have tried several different dopants
to try and lower the sintering temperature of BMT while maintaining the microwave dielectric
properties, but if the sintering temperatures are lowered, the Q-factor is also significantly
lowered.

2.2.1.2 Solid State Sintering. Sintering refers to the process of firing and consolidating a body
shaped from powder particles. When powder particles are compacted, they are in contact with
each other at many sites with a large amount of pore space between them. The powders that
have been dried and pressed are traditionally called “green products”. These are heated in a furnace to develop the desired structure and properties. This is sufficient to cause significant atomic diffusion for solid-state sintering or significant diffusion and viscous flow when a liquid phase is present or produced by a chemical reaction. The driving force for sintering is the capillary pressure associated with surface curvature where the particles come into contact. In order to reduce the surface area, atoms diffuse to the points of contact, permitting contact flattening. Contact flattening involves the removal of matter between particle centers by volume diffusion, grain boundary diffusion or diffusion through a liquid phase. These mechanisms must be promoted for the pores to shrink and eventually be eliminated. Coarsening mechanisms such as surface diffusion or evaporation-condensation only redistribute matter around surfaces and must be minimized. There are three major stages of sintering. In the initial stage, there is surface smoothing of particles, neck growth, rounding of interconnected open pores. Diffusion of active, segregated dopants can occur and the porosity decreases by <12%. In the intermediate state, there is a formation of continuous pores channels along many three-grain junctions. As densification proceeds, the pore channels form closed pores at densities greater than 92%. In the final stage of sintering, isolated pores are eliminated by mass transport from the grain boundary to the pore. Pores on a grain boundary can be eliminated by grain boundary or lattice diffusion but pores within grains can only be eliminated by lattice diffusion, which can be a problem since volume diffusion often has a higher activation energy. In most ceramics, the lattice diffusivity of the slower rate-limiting constituent is often too slow for effective annihilation of pores that have become trapped in grains. If high density is desired, it is important that the pores remain attached to the grain boundary. In the initial stage there are a number of competing paths for material to be transported to the neck area. Some of these paths will lead to densification, which
is a shrinkage process that requires the centers of the particles to approach. Other paths will lead
to coarsening, which is the growth of the neck between the particles leading to a reduction in
specific surface area without shrinkage. Generally, densification is desired. These different
paths of diffusion are competitive paths of transport; the mechanism that gives the fastest rate of
neck growth will dominate and cause either densification or coarsening. Surface diffusion is the
general transport mechanism that can produce surface smoothing, particle joining and pore
rounding but there is no volume shrinkage. In materials where vapor pressure is high,
sublimation and vapor transport to surfaces of lower vapor pressure also produce these effects.
In contrast, grain boundary and lattice diffusion produce neck growth and volume shrinkage. At
high temperatures, grain boundary diffusion and lattice diffusion tend to dominate and therefore
promote densification. Firing at a low temperature can cause coarsening without significant
densification and can be useful in making porous materials. Therefore the enhanced solid state
sintering of DR’s requires dopants that enhance volume and grain boundary diffusion at lower
temperatures without significantly enhancing defects responsible for extrinsic defects that effect
Q. Solid state sintering of mixed cation oxides such as those used in DR’s is an ambipolar
diffusion process in which anions and cations must be transported from the source to the sink
without kinetic de-mixing. The rate limiting mechanism will be the slowest moving species
along the fast path. Therefore, sintering depends a great deal on the diffusivities; and additives
which influence diffusion can have an enormous impact on the rate of sintering. Additives can
often be beneficial to achieving high densities. Additives can do a number of things including:
decrease the rate of coarsening in the earlier stages of sintering, increase the rate of densification,
decrease the rate of grain growth or increase pore mobility such that pore-grain boundary
separation does not occur.
2.2.1.3 **Liquid Phase Sintering.** When a wetting liquid is present, bulk viscous flow can cause volume shrinkage. However, commercial ceramics tend to have a small amount of reactive liquid that accelerates the densification rather than facilitate viscous flow. A simple type is a eutectic liquid in which the primary phase is partially soluble is often used in perovskite systems. This can provide a wide variety of results depending on the rate of process, the amount of liquid formed at the sintering temperature, the relative solubility’s of the additive and the principle constituents and the wetting behavior of the liquid. In order for the additive to be successful there must be an appreciable solubility for the principle constituent, a reasonably low viscosity for rapid diffusion kinetics and a contact angle that allows it to wet and penetrate between the constituent particles or grains. When the liquid coats each grain, the material can often be sintered to a higher density at a lower temperature with less of a tendency for exaggerated grain growth. Less than 1 vol% liquid phase is sufficient to coat the grains if the liquid is distributed uniformly and the grain size is about 1 µm. The wetting liquid concentrates at the particle contacts and forms a meniscus; this exerts an effective compressive pressure on the compact. There is a rapid rearrangement of particles into a higher density configuration. This is often where the greatest densification takes place. After the initial rearrangement, further densification takes place as particle contacts flatten under the compressive stress applied to the point contacts by capillary pressure. The contact flattening occurs through dissolution at the particle contacts and transport of the materials towards stress-free interfaces. Since the solubility of the smallest particles is the greatest dissolving grains reprecipitate onto larger grains in a way that leads to flattened boundaries and a more densely packed array of grains. This can lead to appreciable grain growth compared to solid state sintering. In some systems, significant changes can occur during the cooling stage. The solid and liquid compositions change frequently with precipitation.
of the principal phase. As the sample cools below the solidus temperature, crystallization of a 
liquid phase can cause a discontinuous increase in the dihedral angle to a non-wetting 
configuration, since the solid-solid interfacial tensions are usually greater than the solid-liquid 
tensions.

2.2.1.4 Point Defects. When there is a decrease or increase in the Q-factor of BMT ceramics 
many researchers attribute the change to point defects, although very few go into detail about 
specifics. Kim et.al.\(^{(32)}\) looked at defects in BMT doped with BaWO\(_4\). When the W\(^{6+}\) is 
assumed to substitute on Ta\(^{5+}\) position there will be positively charged regions created due to the 
charge difference between the cations. TEM showed that there was only 1:2 ordering, and 
therefore compensation did not arise from negative regions created from 1:1 ordering. It is 
therefore probably compensated by the generation of negatively charged point defects. Kim 
determined the main defect mechanism in pure BMT was the formation of Ba vacancies by 
measuring the electrical conductivity as a function of the partial pressure of oxygen and by 
considering charge neutrality and the mass action law; therefore BMT is a p-type conductor. He 
also determined that the main defect created by the doping of WO\(_3\) would be the formation of 
tantalum and oxygen vacancies.

2.2.1.5 B-Site Cation Ordering. Some research groups\(^{(61-64)}\) believe that long range B-site cation 
ordering can significantly increase the Q-factor of BMT and related ceramics. This conclusion 
comes from the Q-factors and the ordering parameters both increasing when the ceramics are 
annealed for long periods of time. This does suggest that there may be some correlation but 
there have also been many reports that seem to contradict this theory. As mentioned earlier, sol-
gel synthesis can produce ceramics that are almost fully ordered and still have a low Q-factor\(^{(43)}\). 
Ra et. al\(^{(57)}\) produced BMT through a conventional mixed oxide route and by a columbite route
(to produce ceramics that were more highly ordered). The ceramics produced by the columbite route had an ordering parameter that was 20% higher than the BMT produced by the conventional mixed oxide route but the Q-factor only increased by about 8%. Matsumoto et. al.\(^{(60)}\) doped BMT with BaSnO\(_3\) and found that with increased dopant the Q-factor increased but the ordering parameter deceased. And Yoon et. al.\(^{(50)}\) had similar results with BMT doped with BaWO\(_3\). Therefore although it does seem reasonable that the amount of B-site cation ordering may have some influence on the Q-factor it is probably not the dominant factor. There are many other sources of extrinsic defects.
**Figure 1** Microwaves are electromagnetic radiation. They lie in the region of $10^{-1}$ to $10^{-3}$ m.
Figure 2 The microwave spectrum and some of the current applications utilizing microwaves. (65)
Figure 3 Fields in a microwave dielectric resonator in the simplest standing wave mode (TE-O1d): (a) magnetic field, H; (b) electric field, E; (c) variation in Ef and Hz with r at Z=0 with reference to cylindrical coordinates.
Figure 4  Mechanisms of dielectric polarization.
Figure 5 A sample resonance curve that is produced when a DR is coupled to a microwave circuit.
Figure 6 The parameter used describes the signals that the network analyzer is measuring. The dielectric constant is measured by the $S_{21}$ parameter and the Q-factor is measured using the $S_{11}$ parameter.

\begin{align*}
S_{11} &= \frac{b_1}{a_1} \quad \text{Return Loss} \\
S_{21} &= \frac{b_2}{a_1} \quad \text{Insertion Loss} \\
S_{12} &= \frac{b_1}{a_2} \quad \text{Isolation} \\
S_{22} &= \frac{b_2}{a_2} \quad \text{Output Return Loss}
\end{align*}
The Q-factor of a microwave ceramic can be measured by using the single port reflection technique. The DR is placed on a piece of Styrofoam in the center of a silver cavity. Microwaves are sent from the network analyzer into the cavity. The signal bounces around the cavity and exits though the same port as it entered.

**Figure 7** The Q-factor of a microwave ceramic can be measured by using the single port reflection technique. The DR is placed on a piece of Styrofoam in the center of a silver cavity. Microwaves are sent from the network analyzer into the cavity. The signal bounces around the cavity and exits though the same port as it entered.
Figure 8  The Courtney method can be used to measure the dielectric constant of a microwave ceramic. The ceramic is placed between two metal plates. The signal comes from the network analyzer into the cavity and leaves the opposite side back into the network analyzer into the cavity and leaves the opposite side back into the network analyzer.
Figure 9  The perovskite structure of BMT. The barium ions occupy the corner positions of the cube. The magnesium and tantalum ions occupy the body-centered positions. And the oxygen anions occupy the face-centered positions.
3.0 OBJECTIVES

The primary objective of this research was to find a dopant(s) that could lower the sintering temperature of barium magnesium tantalate (BMT) while maintaining or improving the microwave dielectric properties. The sub-tasks that were undertaken to achieve this goal were:

- Investigate additive(s) that would form a meta-stable liquid phase that would promote densification while present and then go into solution.
- Investigate additive(s) that would create oxygen vacancies to try and increase the diffusivity of oxygen. Since oxygen is the largest ion it is most likely rate limiting during sintering, therefore adding oxygen vacancies may help to aid sintering. However none of the diffusion coefficients are known.
4.0 EXPERIMENTAL PROCEDURE

4.1 Synthesis of BMT Powders

High purity (>99+%% BaCO$_3$ (Aldrich), MgO (Aldrich), Ta$_2$O$_5$ (H.C. Stark), ZnO (Aldrich) and Ga$_2$O$_3$ (Puratronic) were used as starting materials. The chemical analysis from the manufactures are given in appendix A. The as-received magnesium oxide was calcined in a covered alumina crucible at 1000°C for three hours with a heating and cooling rate of 10°C/min. At 300°C the MgO was removed from the furnace and weighed into plastic bottles. This calcination rids the powder of any carbonate and hydroxide absorbed from the air, assuring, as much as possible, the correct stoichiometry of MgO during mixing. The MgO was difficult to weigh accurately because of the temperature and the rate at which moisture was adsorbed, both caused the weight to fluctuate. Appropriate amounts of starting powders (BaCO$_3$, calcined MgO, Ta$_2$O$_5$, ZnO and Ga$_2$O$_3$) were milled in a high-density polyethylene bottle for 24 hours. (The powder amounts are shown in table 1.) The powders were milled using 5-mm yttria-stabilized zirconia (YSZ) balls (YSZ; Union Process Inc., Akron, OH.) and acetone as carrier fluid. The mixed slurry was dried on a hotplate (PMC 730 series, Dataplate; PMC Industries, Inc., San Diego CA.) at 60°C with continuous stirring at 200 RPM to prevent segregation by differential settling. The powder was then kept in a drying oven (OV-18A; Blue M, Hawthorne, CA.) at 108°C for at least 12 hours to ensure complete dryness. The mixed powder was ground with a mortal and pestal and then calcined in a platinum crucible in air at 1300°C for 5 hours with a
10° C per minute heating and cooling rates (Lindberg model 51524; Lindberg, Watertown, WI.). The calcined powders were then ball milled again in acetone for 48 hours to break any agglomerates. The mixed slurry was dried on a hotplate at 60° C with continuous stirring at 200 RPM to prevent segregation. The powder was kept in the drying oven at 108° C until further processing. A typical batch produced 30 grams of powder. A schematic of the processing procedure is shown in figure 10 and a detailed processing procedure is given in appendix B.

For the synthesis of the BMT doped with 3% excess MgO; 8mol% ZnO and Ga2O3; 4mol% ZnO and Ga2O3; 8mol% ZnO; and 8mol% Ga2O3, the dopants were added before the first ball milling step. For the BMT doped with ZnGa2O4, the dopant was added after the calcination step but before the second ball milling. The ZnGa2O4 was prepared similarly to BMT. The appropriate amounts of ZnO and Ga2O3 powders were weighed and ball milled in a high-density polyethylene bottle for 24 hours with acetone as the carrier fluid. The powder was dried and kept in the drying oven for 24 hours and then calcined in a platinum crucible with an alumina cover at 1100° C for 24 hours. The powder was then ball milled again for 24 hours in acetone and dried on the hot plate then kept in the drying oven at 108° C until further processing. XRD analysis confirmed ZnGa2O4 (3-1155).

4.2 Powder Processing

The BMT powder was ground with a mortar and pestel and made into pellets using single action uniaxial pressing. Portions of approximately 3.92 grams of sample were weighted and pressed into cylindrical pellets with a diameter of 0.5 inches (1.27 cm) and a height of
approximately 0.3 inches (0.76 cm) using 1,000 psi. The pellets were then cold isostatically pressed at 40,000 psi. (R-4-12-60; HIP) The pellets were enclosed in latex (non-lubricated condoms) to keep them from contacting the oil.

The samples were then sintered in a covered platinum crucible and for various temperatures between 1300°C and 1600°C for times between 0 and 24 hours.

Several samples that were doped with 8mol% ZnO and Ga₂O₃ were cut in half and polished to a 1μm finish using diamond paste. The samples were then thermally etched at 1300°C for 12 minutes. This made it possible to examine the microstructure more carefully.

4.3 Microstructural Characterizations

Phase evolution and microstructure characteristics were examined using x-ray diffraction (XRD, Philips X’PERT systems) and scanning electron microscopy (SEM, Philips XL30FEG).

4.3.1 X-ray Diffraction (XRD) Analysis

The calcined powders and sintered surfaces of the ceramics were characterized by XRD. The scans were done on both sides of the sintered pellets to identify any discrepancies. The JCPDS pattern of BMT (18-176) and of zinc gallate (3-1155) are shown in figure 11. The standard Bragg-Brentano geometry pattern was used to characterize the ceramics. The samples were attached to the stage using double-sided tape. The sample and the detector rotate in such a
way that the incident angle and detecting angle are $\theta$ and $2\theta$ respectively. The tube voltage was 30 kV and the current was 20 mA. CuKα radiation ($\lambda=1.54\text{Å}$) was used with a detector slit of $\frac{1}{2}$ degree and a source slit of 0.3. The materials were scanned using a continuous scan from 5 to 80 degrees ($2\theta$) with a scan speed of 0.040°/sec.

4.3.2 Scanning Electron Microscopy (SEM) Characterization

Before the SEM analysis, the as-sintered sample surfaces were coated with Pd to reduce charging problems. The SEM was operated at 10-15kV. The surface morphology, grain size, and microstructure were examined using the secondary electron imaging and the backscatter imaging. The backscatter image can be helpful in locating/identifying second phases; phases containing heavier elements appear brighter in the backscatter image. The compositions of the ceramics were analyzed using EDX.

4.3.3 Density Measurements

The density of the sintered ceramics was determined using the dimensions and the weight of the samples. The dimensions were measured using a vernier caliper with a 150 mm scale in 0.05 mm subdivisions. The density was then calculated using equation 4-1.

$$\text{Density} = \frac{\text{weight}}{((\delta)(\text{radius})^2(\text{height}))}$$  \hspace{1cm} 4-1
The percent density was then calculated by dividing the calculated density by the theoretical density. It was assumed that the small amount of dopants (ZnGa$_2$O$_4$, MgO,ZnO and Ga$_2$O$_3$) did not affect the BMT theoretical density of 7.628 g/cm$^3$.(26)

4.4 Microwave Dielectric Property Measurements

The microwave dielectric properties were measured using a network analyzer (Hewlett Packard 8722C). A network analyzer measures the reflection and transmission characteristics of devices by applying a known swept signal and measuring the response of the device. The signal that is transmitted or reflected from the device is compared to the incident signal generated by a swept RF source. The built-in synthesized source generates a continuous wave signal of 40 GHz with a frequency resolution of 100 KHz. The source is generated by a voltage controlled oscillator.

4.4.1 Measurement of the Dielectric Constant

The dielectric constant of a material describes its ability to polarize and store electrical charge. It is simply the ratio of the permittivity of the material to the permittivity of a vacuum. The relative dielectric constant (k’) of the sintered ceramics was measured using the Courtney cavity method$^{(9)}$ and the TE-011 mode. The system set-up is shown in figure 12. This is a transmission type measurement where the signal enters through one side of the cavity and is
collected at the other side. This measurement uses the $S_{21}$ parameter; the $S_{21}$ parameter is equal to the electric field of the microwave signal leaving the component output divided by the electric field of the microwave signal entering the component input. A small brass cavity known as a Courtney cavity is utilized; the sample is secured between the top and bottom of the cavity. The resonance frequency in the TE-011 mode (The mode describes the direction of the magnetic fields used.), 3 dB half power bandwidth, and the insertion loss at the resonant frequency were obtained from the network analyzer. Using this information a software program developed by GDK Products, Cazenovia, NY. was used to calculate the dielectric constant. The accuracy of the Courtney method for measuring the dielectric constant is within 1%.\(^9\) A detailed procedure is given in appendix C. Measurements were taken from several, randomly selected pellets from each condition.

4.4.2 Measurement of the Quality Factor

The quality factor, $Q$, is the ability of a dielectric resonator to store microwave energy with minimal signal loss. The $Q$ is inversely proportional to the microwave loss, therefore high $Q$’s are desirable. The unloaded quality factor ($Q_d$) of the sintered ceramics were measured using the single port reflection technique (the signal enters and exits through the same port in the cavity) and a silver plated cylindrical aluminum cavity designed originally by engineers at Trans-Tech Inc., Adamstown, MD. The sample was isolated from the cavity walls with a styrofoam support. This is done to reduce conduction. This set-up is shown in figure 13. The sample is placed on the support and centered in the cavity. The network analyzer then sweeps the frequency range and gives a number of resonant modes. This measurement technique uses the
$S_{11}$ parameter; the $S_{11}$ parameter is the electric field of the microwave signal leaving the component input divided by the electric field of the microwave signal entering the component input. The TE-01 (also referred to as the TE01$\delta$) mode is used for this measurement; this is the lowest frequency mode. There are two ways to confirm that the mode is the TE-01 mode. When the lid of the cavity is opened, the TE-01 mode will not disappear. And when a piece of metal is inserted into the cavity, the TE-01 mode will shift up in frequency. This is because the metal has a higher dielectric constant than the ceramic. When the resonant frequency is located, the system is set to span 1GHz (with the resonant frequency near the center). The system is then calibrated for this frequency range using a Hewlett Packard calibration kit. The return loss ($S_{11}$) is read directly from the network analyzer. The modified return loss (RL') is calculated from the return loss (RL) read from the network analyzer using a program from John Deriso of Trans-Tech, Inc. (This is done to correct for non-perfect coupling.) Even though the RL' can be calculated for any over or under resonant coupling, all the resonant couplings were adjusted to almost critical coupling by moving the probe. Critical coupling is observed when the coupling coefficient is unity. The coupling coefficient is defined as the ratio of the power dissipated in the external circuit to the power dissipated within the shielded cavity. Therefore, if the power dissipated in the external circuit is large in comparison with the power dissipated in the unloaded cavity, the coupling coefficient is large. Coupling is determined using the Smith Chart, a sample smith chart is shown in figure 14. Using the RL' as the half power bandwidth value, the network analyzer calculates the unloaded quality factor. A detailed procedure is given in appendix C. This technique is believed to lead to a $Q_d$ measurement within 2%. ($^{56}$)
Figure 10  The procedure for the synthesis of BMT. When BMT was doped with ZnO, Ga2O3, and MgO, the dopant(s) was added before the first ballmilling. When the BMT was doped with ZnGa2O4, the dopant was added before the second ballmilling.
Figure 11  The JCPDS X-ray diffraction pattern for BMT (18-176).
Figure 12  The JCPDS X-ray diffraction pattern for ZnGa$_2$O$_4$ (3-1155).
Figure 13  Courtney method set-up for the measurement of the dielectric constant.
Figure 14 Single port reflection technique set-up for the measurement of the quality factor.
Figure 15 A smith chart showing critical, over and under coupling.
Table 1  Starting Materials for Various Compositions of BMT

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<tr>
<th>Composition of BMT</th>
<th>BaCO$_3$ (g)</th>
<th>Ta$_2$O$_5$ (g)</th>
<th>MgO (g)</th>
<th>ZnO (g)</th>
<th>Ga$_2$O$_3$ (g)</th>
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<td>0.6288</td>
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5.0 RESULTS

5.1 BMT Doped with 8mol% ZnO and Ga2O3

5.1.1 Powder Development

The calcined powder of BMT doped with 8mol% ZnO and Ga2O3 was examined by x-ray diffraction and scanning electron microscopy. A typical x-ray diffraction pattern is shown in figure 15. All of the powders showed only BMT peaks with no signs of any second phase or any chemically different material by EDS. A typical SEM micrograph is shown in figure 16. This powder showed no signs of second phase. This does not mean that there is no minor crystalline second phase; it may just be below the detection limit of the XRD (=2 vol%).

5.1.2 Ceramics Sintered at a Constant Heating Rate

BMT was co-doped with 8mol% ZnO and Ga2O3 before the first ball milling in the synthesis. These ceramics were sintered at a constant heating rate of 10°C/min (heating and cooling) to various temperatures between 1300°C and 1600°C. Figure 17 shows a graph of temperature verses density for these ceramics. The density increases up to 1500°C for 0 hours (density = 94.9 % ± 1.0), the sample heated to 1600°C for 0 hours showed essentially the same
density (density = 94.7\% \pm 0.3). As can be seen from the graph, the largest density increase is between the temperatures of 1400\(^\circ\)C and 1450\(^\circ\)C where the density increases from 79.6\% \pm 3.4 to 92.9\% \pm 1.3, suggesting this is a viable temperature range for isothermal sintering without extended sintering times. Figure 18 shows a graph of the same ceramics and conditions with Q\(^*\)f versus temperature. The Q-factors start very low (Q\(^*\)f = 24,108 GHz \pm 3,014) and then quickly increase between 1400\(^\circ\)C and 1450\(^\circ\)C (37,013 GHz \pm 10,404 to 118,283 GHz \pm 5,855); this is much steeper than the graph of density versus temperature. The large increase in the Q-factor corresponds to the large increase in density from \(\approx\)80\% to over 94\%, probably when the pores begin to close off. Therefore, a large increase in the Q-factor seems to correlate with densities above which pore closure is expected. The Q\(^*\)f reaches a maximum of 122,247 GHz \pm 6,629 GHz at 1500\(^\circ\)C but is only increased slightly from 118,283 GHz \pm 5,855 at 1450\(^\circ\)C, so there is no statistical difference.

Selected pellets from each condition were examined using x-ray diffraction and scanning electron microscopy. Several SEM micrographs of the sintered surfaces are shown in figure 19. The microstructure of the sintered surface of ceramics sintered at 1300\(^\circ\)C and 1350\(^\circ\)C were similar. The grain size was less than 0.5 \(\mu\)m and many fine pores were observed, consistent with the measured density. These ceramics showed no signs of any second phase by SEM but it was hard to tell because of the large amount of porosity. When the temperature was increased to 1400\(^\circ\)C the microstructure was less porous, the grains were approximately 1 \(\mu\)m and there were some signs of second phase. The second phase consisted of small crystals that contain mostly Zn and Ga. These crystals could be imaged in backscatter and were very small (<0.5 \(\mu\)m) and difficult to identify. As the temperature was increased to 1450\(^\circ\)C and 1500\(^\circ\)C the crystals of second phase became much more evident (larger) and there seemed to be an increased number of
them. And when the temperature reached 1600°C the crystals were easily identified and some of
them were very large (>20 µm). (There was only one pellet sintered at 1600°C and therefore it
was not included on the graphs showing the density and Q-factor.)

The x-ray analysis showed only BMT peaks (82-0371) for most of the ceramics. In some
of the patterns there were small second phase peaks or what must be from a second phase. These
peaks were identified as either ZnGa$_2$O$_4$ (71-0843) or Ba$_3$Ta$_5$O$_{15}$ (83-0713). The ZnGa$_2$O$_4$ could
have formed from a reaction between the ZnO and Ga$_2$O$_3$ therefore the zinc and gallium rich
second phase is thought to be zinc gallate. The Ba-Ta phases can form from MgO evaporation.
A typical pattern is shown in figure 20 with the second phases marked.

5.1.3 Ceramics Sintered Isothermally at 1350°C

BMT doped with 8mol% ZnO and Ga$_2$O$_3$ was isothermally sintered at 1350°C from no
dwell time up to 24 hours. This temperature was chosen to allow the kinetics of isothermal
sintering to be investigated. Sintering to full density would probably be done above 1400°C.
The commonly used semi-log plot is shown in figure 21 and has an $R^2$ value of 0.9737. From 0
time to 2 hours the density increases linearly and then begins to level off. Between 2 hours and 6
hours the density is statistically the same at 86.5% ± 5.5 and 90.8% ± 1.8. For 6, 10 and 24
hours there is little difference in the density (90.8% ± 1.8, 89.8% ± 1.8, and 93.8% ± 1.2). The
Q*$f$ verses time is shown in figure 22. The Q-factor reaches a maximum at 10 hours with Q*$f$
equal to 142,907 GHz ± 2,770 GHz. The large change in the Q-factor again seems to correlate
with reaching relative densities above 90%. This further supports the hypothesis that isolated porosity is a major requirement for high Q-factors and low dielectric losses.

Selected pellets from each condition were examined using x-ray diffraction and scanning electron microscopy. Several SEM micrographs of the sintered surfaces are shown in figure 23. Micrographs for samples sintered at 0 time and 1 hour were very porous with the average grain size less than 1 \( \mu \text{m} \); there were no signs of any second phases. At 2 hours the microstructure looked less porous with an average grain size of about 1 \( \mu \text{m} \) and there were some second phases of Zn-Ga rich crystals as mentioned in the last section. At sintering times of 6, 10 and 24 hours the micrographs looked very similar. There were more Zn-Ga crystals and they were larger. In samples sintered for 2 hours the Zn-Ga crystals were less than 0.5 \( \mu \text{m} \), and in the samples sintered for the longer times the crystals were 2 \( \mu \text{m} \) or larger. Figure 24 shows the fracture surface of a pellet sintered at 1350\(^\circ\)C for 24 hours. The fracture goes through the BMT grains (intergranular) and around the second phase crystals (transgranular). The second phase crystals appear to be growing out of the BMT solid and there are no pores around crystals. This suggests that the zinc gallate crystals nucleate and grow in the solid-state.

The x-ray analysis showed only BMT peaks for 0 time and 1 hour. In all of the longer times there were some very small second phase peaks. These peaks were identified as either ZnGa\(_2\)O\(_4\) or Ba\(_3\)Ta\(_5\)O\(_{15}\). As mentioned earlier, ZnO and Ga\(_2\)O\(_3\) can react to form ZnGa\(_2\)O\(_4\) and MgO can evaporate to form a Ba-Ta rich phase. The intensities of these peaks varied slightly from pellet to pellet even when sintered at the same condition. A typical pattern is shown in figure 25. The second phase peaks are marked appropriately. The pellet shown was sintered at 1350\(^\circ\)C for 24 hours; both the top and bottom of the sintered surfaces of the pellet are similar.
5.1.4 Ceramics Sintered at Various Temperatures for 2 hours

BMT co-doped with 8mol% ZnO and Ga$_2$O$_3$ was sintered at various temperatures between 1300°C and 1600°C for 2 hours. Two hours was chosen because it is a desirable sintering time for manufacture. This also allowed the samples to equilibrate more than in the constant heating rate work. The graph of temperature verses density is shown in figure 26. The density increases linearly up to 1400°C to a density of 94.5% ± 0.5 and then drops slightly at 1450°C to 93.9% ± 1.2. As the temperature increases to 1500°C and 1550°C the density remains about constant at 93.2% ± 0.7 and 93.4% ± 0.7. At 1600°C the density drops slightly to 89.6% ± 3.2. Compared to the constant heating rate work the temperature interval in which densification accelerates is shifted to lower temperatures as expected. The plot for Q*f verses temperature is shown in figure 27. The Q-factor increases rapidly from 1300°C to 1350°C (although the error is quite large for the pellets sintered at 1350°C), and then increases much slower between 1350°C and 1500°C. Between 1550°C and 1600°C there is a slight drop in the Q-factor from 139,565 GHz ± 8,762 to 97,737 ± GHz 6,756. Therefore the relationship of temperature with the Q-factor and temperature with density are similar. The results are still consistent with high Q-factors requiring closed pores.

Selected pellets from each condition were examined using x-ray diffraction and scanning electron microscopy. The micrographs from these ceramics are very similar to those that were sintered for 0 time at the same temperatures. The second phase of Ga-Zn rich crystals did appear at 1400°C instead of 1450°C where they first became evident in the samples sintered at 0 time. This is consistent with what would be expected if the system is equilibrating during the sintering process and Ga-Zn crystals are nucleating and growing. Another difference was in the samples
sintered at 1600°C. The sintered surface of these ceramics varied from one part to another and also between the different sides of the pellet. Several of these micrographs are shown in figure 28. Some micrographs showed very dense BMT with no second phases (b), some showed very large (some >10µm) second phase crystals (Zn-Ga) (f), and some showed much the same microstructure as was seen previously (e). There was no evidence of any Ba-Ta phases in these samples. There also seemed to be some cracking in the parts of the microstructure with large amounts of the crystals.

XRD analysis showed similar results to those obtained when the sintering time was 0 hours. The small second phase peaks did occur earlier, at 1350°C. A typical pattern is shown in figure 29. The samples sintered at 1600°C for 2 hours were different than the one sintered for 0 time at the same temperature. The pattern showed very large second phase peaks of ZnGa₂O₄ on the bottom and very small peaks on the top. A sample pattern is shown in figure 30. This suggests that the Zn-Ga rich crystals grow and may segregate to certain surfaces, due to non-wetting behavior. This will be discussed later.

5.2 Comparing the Concentration of ZnO and Ga2O3 on the Sintering and Microwave Dielectric Properties of BMT

BMT was co-doped with 4mol% and 8mol% ZnO and Ga₂O₃. Both were added before the first ball milling step in the synthesis. These ceramics were sintered at 1400°C for 2 hours and 1500°C for 2 hours. At both of these temperatures the densities were slightly higher in the BMT doped with 4mol% ZnO and Ga₂O₃. At a sintering temperature of 1400°C for 2 hours the
densities were 94.5% ± 0.5 and 95.2% ± 0.9; and at a sintering temperature of 1500°C for 2 hours the densities were 93.2% ± 0.73 and 96.1% ± 0.7. Although the density was higher in the BMT doped with 4mol% ZnO and Ga₂O₃ the Q*f values were higher for the BMT doped with 8mol% ZnO and Ga₂O₃ at both sintering temperatures. The Q*f values at a sintering temperature of 1400°C for 2 hours were 129,041 GHz ± 4,614 GHz and 126,102 GHz ± 5,141 GHz; and at a sintering temperature of 1500°C for 2 hours the Q*f ‘s were 151,595 GHz ± 1,918 GHz and 130,936 GHz ± 4,526 GHz. A graph comparing the Q-factors at both concentrations and sintering temperatures are shown in figure 31. This shows that the difference at 1400°C was not significant but at 1500°C the difference is more pronounced.

Selected pellets from each of these conditions were examined using x-ray diffraction and scanning electron microscopy. SEM micrographs showed second phases of crystals that were rich in Zn and Ga for both compositions and sintering conditions. These crystals are the same as the Zn-Ga rich crystals discussed earlier. Typical SEM micrographs from each condition are shown in figure 32. The Zn-Ga rich crystals appear darker than the “normal” BMT grains in the images obtained from the backscattered electrons. The BMT doped with 4mol% ZnO and Ga₂O₃ and sintered at 1500°C for 2 hours also showed long, cylindrical second phases that are rich in barium and tantalum. These second phases are only seen on the top of the pellets. The XRD confirms these findings showing large second phase peaks on the tops of these pellets. A typical XRD pattern is shown in figure 33. These second phase peaks also appear on the bottom of the pellets but are very small. These peaks were identified as Ba₃Ta₅O₁₅ (83-0713), which would be consistent with MgO evaporation. Not all of the intensities match exactly with the JCPDS patterns but this can happen when the grains/crystals are not randomly oriented when they grow on the surface. The surfaces of a pellet sintered at 1500°C for 2 hours were removed to try and
remove the Ba-Ta phases. These XRD patterns are shown in figure 34; the second phase peaks disappeared almost completely. The removal of these phases lowered the Q-factor from 27,834 to 25,917, although this could be due to surface grinding. The XRD patterns for the pellets sintered at 1400°C for 2 hours showed small second phase peaks similar to the bottom of the pellets sintered at 1500°C for 2 hours. A typical XRD pattern for BMT doped with 8mol% ZnO and Ga₂O₃ is shown in figure 35. The patterns for both conditions are the same along with both sides of the pellets. The patterns showed all of the BMT peaks and some small second phase peaks. These peaks correspond to Ba₃Ta₅O₁₅ (83-0713), which can form from MgO evaporation and ZnGa₂O₄ (71-0843), which can form through the reaction of ZnO and Ga₂O₃. The intensities of the small second phase peaks varied slightly from batch to batch, almost disappearing in some cases. The results are consistent with more MgO evaporation and barium tantalate formation on the surface of 4mol% ZnO and Ga₂O₃. The effect of the sintering additive on this is unclear but could be due to Zn dissolution.

5.3 Comparing the Effect of Several Similar Dopants (ZnO and Ga₂O₃; Ga₂O₃; ZnO and ZnGa₂O₄) on BMT Sintered at 1450°C for 2hours

BMT was doped with several variations of zinc oxide and gallium oxide to determine the effect of each dopant. BMT was doped with 8mol% of each of the following: ZnO and Ga₂O₃; Ga₂O₃; ZnO; and ZnGa₂O₄. The ZnO and Ga₂O₃ were added before the first ball milling of the BMT; the ZnGa₂O₄ was added after the BMT calcination step in the synthesis. Figure 36 shows a graph comparing the Q* values and the densities. It is clear that the BMT doped with ZnO
and Ga$_2$O$_3$ has the highest Q$^*_f$ at 135,164 GHz ± 17,825 GHz although it also had the lowest density at 93.9% ± 1.2. The BMT doped with Ga$_2$O$_3$ and with ZnGa$_2$O$_4$ had Q$^*_f$'s of 84,391 GHz ± 4,313 GHz and 93,743 ± 33,520 GHz even though they both had average densities over 95%. (BMT doped with Ga$_2$O$_3$ had a density of 95.2% ± 0.6 and BMT doped with ZnGa$_2$O$_4$ had a density of 95.2 ± 0.6.)

SEM micrographs of the BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1450°C for 2 hours showed Zn-Ga rich second phase crystals as described before. This is shown in figure 37. These crystals appear to be zinc gallate that can form from a reaction between the ZnO and Ga$_2$O$_3$. As mentioned earlier, when the fracture surface is examined, the fracture goes through the BMT grains and around the Zn-Ga rich crystals. Another important finding was that there are no pores around the second phase crystals. As would be expected given the density results the microstructures show more porosity that the ceramics with the other dopants sintered at the same conditions. The XRD patterns looked similar to samples sintered at 1500°C for 2 hours and shown in figure 29. The patterns show all of the BMT peaks and small second phase peaks which are indicated. The small second phase peaks correspond to zinc gallate (38-1240).

SEM micrographs of BMT doped with 8mol% Ga$_2$O$_3$ showed plate-like second phase crystals. These appear darker in the backscatter image and can be seen clearly in figure 38. When a fracture surface was examined using SEM, the fracture went around the Ga-rich second phase and through the BMT grain; there was no surrounding porosity. According to EDS, these are rich in gallium. The XRD of both sides of the pellets are shown in figure 39. The patterns show all of the BMT peaks and a few very small second phase peaks. These peaks are slightly larger on the top of the pellets. These small second phase peaks correspond to gallium oxide (38-1240).
SEM micrographs of BMT doped with 8mol% \( \text{ZnGa}_2\text{O}_4 \) and sintered at 1450\(^\circ\)C for 2 hours is shown in figure 40. This micrograph shows two different second phases: long, cylindrical second phases that are rich in Ba and Ta; and groups of crystals that are rich in Zn and Ga. These crystals appear darker when viewing the backscatter image. The Zn-Ga rich crystals appear to be the same crystals that are seen on BMT doped with ZnO and \( \text{Ga}_2\text{O}_3 \). When the fracture surface was examined it looked like the fracture surface of the BMT doped with 8mol% ZnO and \( \text{Ga}_2\text{O}_3 \). The fracture went through the BMT grains and around the Zn-Ga rich second phases; there was no surrounding porosity. There were also no signs of any Ba-Ta rich phases on the interior surfaces, confirming that these phases were formed by the evaporation of MgO. The XRD patterns of both sides are shown in figure 41. All of the BMT peaks are present and both sides show large second phase peaks; the peaks on the top are larger. The second phase peaks seem correspond to two different second phases, zinc gallate and barium tantalate (\( \text{Ba}_3\text{Ta}_5\text{O}_{15} \)). The peak intensities do not match exactly but this can be caused when the crystal are not oriented randomly when they grow on the surface.

5.4 Overall Effect of Density on the Q-factor

Figure 42 shows a plot of \( Q* f_0 \) verses density for all of the pellets produced. With the exception of series four there is a clear increase in the Q-factor when the density is above about 90% of the theoretical density. This is most likely when the pores of the ceramics are all closed off preventing water moisture from getting into the samples. Water strongly absorbs microwaves, which can significantly increase the losses.
5.5 Effect of Additives and Process Condition on Ordering

In all the XRD patterns, the largest superlattice peak has been marked with an asterisk. These peaks are very small compared with the largest BMT peak indicating that there is very little ordering occurring in the samples or that the antiphase domains are too small to be detected by XRD.

5.6 The Effect of Additives and Processing Condition on the Dielectric Constant

The dielectric constant was measured on a couple of pellets from each condition to be sure that it was not being adversely affected. None of the dopants severely (<24) decreased the dielectric constant of BMT; BMT normally has a dielectric constant of approximately 25.
Figure 16  X-ray diffraction pattern of calcined BMT powder doped with 8mol% ZnO and Ga2O3. This pattern shows only BMT peaks and no superlattice peaks.
Figure 17  SEM micrograph of calcined BMT powder doped with 8mol% ZnO and Ga2O3. There is no evidence of any second phases.
**Figure 18** The density Vs sintering temperature for BMT co-doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at a constant heating rate of 10°C/min.
Figure 19 The $Q^*f$ vs sintering temperature for BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at a constant heating rate of 10°C/min. The largest increase in Q-factor was between the temperatures of 1400°C and 1450°C; this corresponds to densities where the pores would be closing off. (The resonant frequency for these pellets was approximately 4.9 GHz.)
Figure 20  Selected SEM micrographs of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at a constant heating rate of 10°C/min. Some of the zinc and gallium rich second phase has been circled. The only pellet shown that did not show any evidence of Zn and Ga rich second phase is shown in figure (a) and was sintered at 1300°C for 0 hours. The pellets were also sintered at 1400°C for 0 hours (b), 1450°C for 0 hours (c), 1500°C for 0 hours (d), and 1600°C for 0 hours (e).
Figure 21  An XRD pattern of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at a constant heating rate of 10°C per minute increasing to 1500°C and then immediately decreasing at 10°C per minute to room temperature.
Figure 6 Density Vs. Sintering Time for BMT Doped with 8mol% ZnO and Ga2O3 and Sintered at 1350°C

\[ y = 3.7433 \ln(x) + 82.5 \]

\[ R^2 = 0.9737 \]

Figure 22 The density vs. sintering time for BMT doped with 8mol% ZnO and Ga2O3 and sintered at 1350°C
Figure 23 The $Q^*f$ vs. sintering time for BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1350°C. (The resonant frequency for these pellets was approximately 4.9 GHz.)
Figure 24  Selected SEM micrographs of sintered surfaces of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered isothermally at 1350°C. Some of the zinc and gallium rich second phase has been circled. The pellet shown in micrograph (a) has been sintered at 1350°C for 0 hours and shows no evidence of any Zn and Ga rich second phase. The pellet shown in micrograph (b) has been sintered at 1350°C for 2 hours, (c) has been sintered at 1350°C for 6 hours. Micrograph (c) was taken using the backscattered electrons and therefore the Zn and Ga rich crystals appear darker than the BMT grains. And the pellet shown in micrograph (d) has been sintered at 1350°C for 24 hours.
Figure 25  A SEM micrograph of a fracture surface of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1350ºC for 1 hour. Some of the Zn and Ga rich second phase has been circled.
Figure 26 A typical pattern from BMT doped with 8mol% ZnO and Ga$_2$O$_3$. This particular pattern came from a sample sintered at 1350°C/24hr. The circles peak has been identified as the start of barium tantalite. All other peaks have been identified as BMT peaks.
Figure 27  A graph of Density vs. Sintering Temperature (time=2 hrs) for BMT Doped with 8mol% ZnO and Ga$_2$O$_3$. 
Figure 28 The Q*\textit{f} vs. sintering temperature (time=2 hrs) for BMT doped with 8mol\% ZnO and Ga\textsubscript{2}O\textsubscript{3}. (The resonant frequency was approximately 4.9 GHz.)
Figure 29  Selected SEM micrographs of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at various temperatures for 2 hours. Some of the zinc and gallium rich second phase has been circled. Micrograph a) has been sintered for 1450°C for 2 hours and is imaged by the backscattered electrons. Micrographs b) through f) have been sintered for 1600°C for 2 hours.
Figure 30  An XRD pattern of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1500°C for 2 hours. The largest superlattice peak has been marked with an arrow.
Figure 31  An XRD pattern of BMT doped with 8mol% ZnO and Ga₂O₃ and sintered at 1600°C for 2 hours. This was taken on the top of the pellet; the bottom has the same peaks with the second phase peaks smaller. The zinc gallate peaks have been marked with an asterisk and the largest super lattice peak has been marked with an arrow.
Comparing the Concentration of ZnO and Ga2O3 Doped BMT

Figure 32  A Comparison of the $Q^* f$ of 4mol% and 8mol% ZnO and Ga$_2$O$_3$ doped BMT. Although the BMT doped with 4mol% had slightly higher densities, the Q-factors were both higher in the BMT doped with 8mol% ZnO and Ga$_2$O$_3$. 
Figure 33  Selected SEM micrographs of BMT doped with 4mol% (a) and b)) and 8mol% (c) and d)) ZnO and Ga$_2$O$_3$. Micrographs a) and c) have been sintered at 1400°C for 2 hours and micrographs b) and d) have been sintered at 1500°C for 2 hours. Some of the zinc and gallium rich second phase has been circled and the barium and tantalum rich second phase is indicated by a rectangle.
Figure 34  XRD pattern of the top of a sintered pellet of BMT doped with 4mol% ZnO and Ga$_2$O$_3$. The barium tantalate second phase peaks have been circled and the largest superlattice peak has been marked with an arrow.
Figure 35 An XRD pattern of the bottom of a pellet doped with 4mol% ZnO and Ga$_2$O$_3$ and sintered at 1500°C for 2 hours.
Figure 36 An XRD of the top of a pellet of BMT doped with 4mol% ZnO and Ga$_2$O$_3$ and sintered at 1500°C for 2 hours with the outer surface removed.
Figure 37 An XRD of the bottom of a pellet of BMT doped with 4mol% ZnO and Ga$_2$O$_3$ and sintered at 1500$^\circ$C for 2 hours with the outer surface removed.
Figure 38  An XRD pattern of a BMT pellet doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1400°C for 2 hours.
**Figure 39** A Comparison of the $Q^f$ of BMT doped with combinations of ZnO and Ga$_2$O$_3$. 
Figure 40  SEM micrograph of BMT doped with 8mol% ZnO and Ga$_2$O$_3$ and sintered at 1450$^\circ$C for 2 hours. Some of the zinc and gallium rich second phase has been circled.
Figure 41  SEM micrograph of BMT doped with 8mol% Ga$_2$O$_3$ and sintered at 1450°C for 2 hours. Some of the gallium rich second phase has been circled.
Figure 42 An XRD pattern of a BMT pellet (top) doped with 8mol% Ga$_2$O$_3$ and sintered at 1450°C for 2 hours. The largest superlattice peak has been marked with an arrow.
Figure 43  An XRD pattern of a BMT pellet (bottom) doped with 8mol% Ga$_2$O$_3$ and sintered at 1450°C for 2 hours. The largest superlattice peak has been marked with an arrow.
Figure 44  SEM micrograph of BMT doped with 8mol% ZnGa$_2$O$_4$ and sintered at 1450°C for 2 hours. Some of the second phase rich in gallium and zinc is circled and some of the second phase rich in barium and tantalum is indicated by a rectangle.
Figure 45  An XRD pattern of BMT doped with 8mol% ZnGa_2O_4 and sintered at 1450°C for 2 hours.
Figure 46  XRD pattern of the bottom of a pellet of BMT doped with 8mol% ZnGa$_2$O$_4$ and sintered at 1450°C for 2 hours. The zinc gallate second phase peaks have been marked with an asterisk and the largest superlattice peak has been marked with an arrow.
6.0 DISCUSSION

6.1 Phase Evolution

6.1.1 BMT Doped with ZnO & Ga2O3 and ZnGa2O4

When the micrographs and XRD patterns of the BMT co-doped with ZnO and Ga2O3 and ZnGa2O4 are compared, the co-doped additives appear to react to form ZnGa2O4 (zinc gallate). This is evident when the second phase crystals are compared. The second phase peaks on the XRD patterns confirm zinc gallate crystals in both instances. This suggests that there is no reason to react the zinc gallate before it is added to the BMT as a sintering aid; and the BMT doped with ZnO and Ga2O3 actually had slightly lower losses. The only difference is the Ba-Ta rich second phases that appear in some of the BMT samples doped with pre-made zinc gallate. This can be caused by the evaporation of MgO during sintering. In the BMT doped with zinc and gallium, the zinc may react with the Ba-Ta rich phases to form BZT. There were no signs of any BZT on any micrographs or XRD patterns but it may be below the detection limit of the instrumentation. The XRD patterns and the SEM micrographs of the samples doped with 4mol% and 8mol% ZnO and Ga2O3 were similar. The micrographs of the samples doped with 8mol% seem to have more second phase crystals but since a quantitative study was not done it cannot be said for certain.
When BMT is doped with zinc and gallium they react during the heating to form zinc gallate. Therefore, after the heating we are assuming that the BMT doped with zinc oxide and gallium oxide and the BMT doped with zinc gallate will react similarly. There was no evidence of any zinc gallate in any micrographs or x-ray diffraction patterns of the calcined powder. There is probably some present it is just too small to be detected. The pre-made zinc gallate was calcined and formed at 1100°C for 24 hours and the BMT doped with gallium and zinc was calcined at 1300°C for 5 hours suggesting that at least some zinc gallate was already formed. There were zinc gallate second phase crystals in most of the micrographs and x-ray diffraction patterns of the sintered surfaces. These second phases grew and became more evident as the temperature and time was increased. On the samples sintered at 1300°C and 1350°C over the first hour there were no signs of any of this second phase on the sintered surface. However, when the fracture surface of a pellet sintered at 1350°C for 1 hour was examined, there were small amounts of second phase present on the interior of the pellet. When the fracture surface of this and others were examined, the fracture went through the BMT grains and around the zinc gallate crystals. Another important observation was that there were no pores around the second phases. This suggests that the phase was formed in the solid state, and that the reaction by which zinc gallate is formed is slow compared with sintering.

6.1.2 BMT Doped with Ga₂O₃, ZnO and MgO

When BMT was doped with Ga₂O₃ (8mol%), ZnO (8mol%), and MgO (3% excess), the solubility of each was exceeded. Therefore, each had large second phase peaks on the XRD patterns and evident second phases on the SEM micrographs. These dopants were added to try
and increase the amount of oxygen vacancies, which is thought to increase mass transport because it is the largest ion and therefore the slowest moving. However it is unclear how much of the dopant actually went into solid solution.

6.2 Densification Behavior

6.2.1 BMT Doped with ZnO & Ga₂O₃ and ZnGa₂O₄

Densification is significantly enhanced by the addition of zinc oxide & gallium oxide and the addition of zinc gallate. The zinc oxide and gallium oxide probably form zinc gallate at the calcination or sintering temperature and then it is coarsening during the sintering. The enhanced sintering could be due to a meta-stable liquid phase that is rich in zinc and gallium. When the system approaches equilibrium, the liquid then forms zinc gallate crystals and there was no residual liquid detected. The materials will have reached equilibrium when the volume fraction of zinc gallate remains constant.

Another possible explanation for the increase in sinterability is that some of the zinc and/or gallium go into solid solution to create barium, tantalum and/or oxygen vacancies. Kim^{[32]} determined that the main defect mechanism in pure BMT was the formation of Ba vacancies and when BMT is doped with WO₃, the main defect mechanism is the formation of tantalum and oxygen vacancies. However, there as no evidence presented to support these statements. The creation of these vacancies may help to increase diffusivity of the barium and
oxygen ions (the largest ions) and therefore increase the sinterability. It is assumed that no interstitials will be formed and due to the size of the ions it is assumed that the gallium will go onto tantalium sites and the zinc will go onto magnesium sites. Table 6.1 shows the ionic radii of the ions. There are many possible defect equations. Given the charge and ionic radii a tentative equation is given below. This creates barium and oxygen vacancies.

\[ \text{Zn}^{2+} + 2\text{Ga}^{3+} + 2\text{O}_2 = \text{Zn}_{\text{Mg}}^{x} + 2\text{Ga}_{\text{Ta}}^{x} + 4\text{O}_{\text{Ba}}^{x} + 3\text{V}_{\text{Ba}}^{x} + 5\text{V}_{\text{O}}^{x} \]

As will be discussed in the following sections, this is probably not the primary mechanism for the increase in sinterability.

6.2.2 BMT Doped with Ga\textsubscript{2}O\textsubscript{3}, ZnO and MgO

The BMT doped with 8mol\% Ga\textsubscript{2}O\textsubscript{3} could be sintered to ≈95% density at 1450°C for 2 hours. It is not clear what is increasing the sinterability. Some of the gallium could be going into solid solution and creating barium and/or oxygen vacancies. The following defect equation is assuming that there are no interstitials are formed; it is also assuming that the gallium will go onto tantalium sites and the oxygen will go onto to oxygen sites.

\[ 2\text{Ga}^{3+} + 3/2\text{O}_2 = 2\text{Ga}_{\text{Ta}}^{x} + 3\text{O}_{\text{O}}^{x} + 3\text{V}_{\text{Ba}}^{x} + 6\text{V}_{\text{O}}^{x} + \text{V}_{\text{Mg}}^{x} \]

Other possible substitutions include,

\[ 3\text{Ga} + 9/2\text{O}_2 = 2\text{Ga}_{\text{Ta}}^{x} + \text{Ga}_{\text{Mg}}^{x} + 3\text{V}_{\text{Ba}}^{x} + 9/2\text{V}_{\text{O}}^{x} + 9/2\text{O}_2 \]

\[ 3\text{Ga} + 3\text{O}_2 = \text{Ga}_{\text{Ta}}^{x} + 2\text{Ga}_{\text{Mg}}^{x} + 3\text{V}_{\text{Ba}}^{x} + 3\text{V}_{\text{O}}^{x} + 6\text{O}_2 \]

all create oxygen vacancies. Alternatively, an unstable gallium rich liquid could form and enhance sintered and then precipitate Ga\textsubscript{2}O\textsubscript{3} as equilibrium is approached.
The BMT doped with 8mol% ZnO and the 3% excess MgO did not sinter at 1600°C for 2 hours. If some of the zinc and/or magnesium could have gone into solid solution it would have created oxygen and barium vacancies, which would be expected to increase the sinterability. It is assumed that there will be no interstitials formed and that the magnesium and zinc will go onto magnesium sites due to the size of the ions.

\[
\text{ZnO} = \text{Zn}_x\text{Mg}^x + \text{O}_x^+ + 3\text{V}_{\text{Ba}}'' + 8\text{V}_O^{**} + 2\text{V}_{\text{Ta}}''''
\]

\[
3\text{Zn} + \frac{3}{2}\text{O}_2 = \text{Zn}_x\text{Mg} + 2\text{Zn}_{\text{Ta}}'''' + 3\text{V}_{\text{Ba}}'' + 6\text{V}_O^{**} + 3\text{O}_O
\]

\[
\text{MgO} = \text{Mg}_x\text{Mg}^x + \text{O}_x^+ + 3\text{V}_{\text{Ba}}'' + 8\text{V}_O^{**} + 2\text{V}_{\text{Ta}}''''
\]

\[
3\text{Mg} + \frac{3}{2}\text{O}_2 = \text{Mg}_x\text{Mg} + 2\text{Mg}_{\text{Ta}}'''' + 3\text{V}_{\text{Ba}}'' + 6\text{V}_O^{**} + 3\text{O}_O
\]

Again the oxygen vacancies are only created. The extent to which this is possible will depend on the limit of stability which is presently not known. The sinterability of BMT was not increased by the addition of ZnO or MgO. Since these addites would also create O and Ba vacancies but would not be expected to form a liquid phase it is thought that a meta-stable liquid phase is the most likely the cause of the enhanced sintering in the gallium containing materials. As equilibrium is established after sintering is complete, zinc gallate or gallium oxide would continue to form by nucleation and growth in dense material. Therefore, porosity is not preferentially located around the second phases. This study did not positively identify the presence of a meta-stable phase and the explanation of the sintering behavior emails a hypothesis requiring a more detailed microstructure study.
6.3 Q-factors

6.3.1 BMT Doped with ZnO & Ga$_2$O$_3$ and ZnGa$_2$O$_4$

For the case of BMT samples doped with ZnO & Ga$_2$O$_3$, and ZnGa$_2$O$_4$, as long as the ceramics got to densities where the pores are closed, the Q-factors were good. This suggests that the second phase crystals do not significantly increase the losses; the losses are measured at relatively high frequency where the interfacial relaxation mechanisms may not significantly contribute to the losses.\(^{(7)}\) This argument would suggest that Q may decrease at lower frequencies when interfacial effects become important.

6.3.2 BMT Doped with each of the following Ga$_2$O$_3$, ZnO and MgO

Although the BMT samples doped with Ga$_2$O$_3$ had high densities the Q-factors were not as high as the samples doped with zinc gallate. As was mentioned earlier, interfacial relaxation does not usually contribute significantly to the losses at these relatively high frequencies. This suggests that the second phases that are formed may strongly absorb microwaves, which could significantly increase the losses. It is not known how strongly gallium oxide absorbs microwaves.

The BMT samples doped with ZnO and MgO did not get dense enough for the pores of the sample to close. As discussed earlier, in order to achieve a high Q-factor the sample must have closed pores. Before the pores are closed off, water moisture is able to get into the sample.
Since water moisture is a strong absorber of microwaves, this significantly increases the losses of the ceramic.

The implication of this study is that multiphase DR ceramics can be manufactured with high Q values. The processing of high purity materials by high temperature sintering and prolonged annealing may not be necessary for all microwave ceramics. Therefore sintering aids and conventional sintering practice are a viable way of lowering processing temperatures even if they create second phases, as long as these phases do not strongly absorb microwaves. At present there is not way of predicting the effects of sintering aid on Q.
Table 2  Ionic Radii of BMT and Other Relevant Ions

<table>
<thead>
<tr>
<th>Ion</th>
<th>Ionic Radius in Å with a coordination number of 6</th>
</tr>
</thead>
<tbody>
<tr>
<td>Ba$^{2+}$</td>
<td>1.35</td>
</tr>
<tr>
<td>Mg$^{2+}$</td>
<td>0.72</td>
</tr>
<tr>
<td>O$^{2-}$</td>
<td>1.40</td>
</tr>
<tr>
<td>Ta$^{5+}$</td>
<td>0.64</td>
</tr>
<tr>
<td>Ga$^{3+}$</td>
<td>0.62</td>
</tr>
<tr>
<td>Zn$^{2+}$</td>
<td>0.74</td>
</tr>
</tbody>
</table>
BMT can be sintered to high densities at 1450°C or lower for 2 hours by the addition of zinc gallate without significantly lowering the Q-factor. The dopant can be added as zinc gallate or as zinc oxide and gallium oxide.

Zinc gallate may lower the sintering temperature of BMT through a meta-stable transient liquid phase.

When BMT is doped with zinc gallate or zinc oxide and gallium oxide, second phase crystals are formed at either the calcination temperature or the sintering temperature and then coarsen during continued sintering.

Multiphase materials can have high Q-factors or low losses at a frequency of around 5 GHz.

BMT doped with 8mol% ZnO and sintered at 1450°C for 2 hours only had a density of approximately 85% but it had a relatively high Q*f of 126,439 GHz. Only one batch of this composition was made so it is not clear if this is significant.

BMT can be sintered to high densities at 1450°C for 2 hours by the addition of gallium oxide, but this significantly lowers the Q-factor of the materials.
8.0 SUGGESTIONS FOR FUTURE WORK

1. A measure the Q-factor of doped BMT ceramics at lower frequencies could determine the effect of the second phases on the loss.

2. A detailed microstructure evaluation needs to be done to determine when the systems reaches equilibrium and if there are any signs of liquid phases sintering. In particular TEM investigation of the chemical composition of second phases collected at triple points must be performed.

3. The temperature coefficient of resonant frequency should be measured to determine if it is severely degraded by the dopants.
APPENDICES
APPENDIX A

Chemical Analysis of Starting Materials

**Barium Carbonate (BaCO₃)**
Purchased from Aldrich  
Product Number: 23710-8  
Formula Weight: 197.35  
Appearance: White Powder  
Impurities:
- Calcium: 0.003%  
- Chloride: 0.0007%  
- Iron: 0.0003%  
- Heavy Metals: <0.001%  
- Sodium: 0.001%

**Gallium (III) Oxide (Ga₂O₃)**
Purchased from Puratronic  
Product Number: 11151  
Formula Weight: 187.44  
Appearance: White Powder  
Impurities: (ppm)
- Al: 0.13  
- As: <0.20  
- Au: <0.50  
- B: 0.08  
- Ca: 0.18  
- Cd: <0.10  
- Co: <0.01  
- Cr: 0.03  
- Cu: <0.05  
- Fe: 0.04  
- Ge: <0.05  
- Hg: <0.10  
- In Binder: Mg 0.19  
- Mn: <0.01  
- Na: 0.97  
- Ni: 0.04  
- P: 0.04  
- Pb: <0.05  
- S: 0.89  
- Si: 4.60  
- Sn: 0.56  
- Te: <0.05  
- Zn: <0.10

**Magnesium Oxide (MgO)**
Purchased from Aldrich  
Product Number: 20371-8  
Formula Weight: 40.31  
Appearance: White Powder
Impurities: (ppm)

- Zn 25
- Ba 0.6
- Al 20
- Zr 0.8
- Na 20
- Li 1.8
- Mn 15
- Mo 3.0
- Ca 7.7
- Ti 0.4
- Pb 7.4
- Ni 0.2
- Cu 4.0
- B 0.1

**Tantalum Pentoxide (Ta₂O₅)**
Purchased from H.C. Starck
Product Number: 99/45904
Formula Weight: 441.89
Appearance: White Powder

Impurities: (ppm)

- Fe <10
- Nb 33
- Ni <10
- Si <10
- Ti <2
- W <5
- Zr <2
- Alkali <30

**Zinc Oxide (ZnO)**
Purchased from Aldrich
Product Number: 20553-2
Formula Weight: 81.37
Appearance: White Powder

Impurities: (ppm)

- Cd 5
- Al 3
- Na 3
- Ca 2
- Fe 2
- Pb 1
- Ba 0.3
- Mg 0.3
APPENDIX B

Procedure for the Synthesis of BMT Ceramics

Sample Preparation

1. Calcine MgO in an alumina crucible with cover at 1000°C for 3 hours. Use a heating and cooling rate of 10°C/min, terminating at 300°C.
2. At 300°C remove the MgO from the furnace.
3. Weigh the MgO immediately and place in clean milling bottles with a few mL of acetone – be sure to write the weight of the MgO on bottle.
4. Weigh stoichiometric amounts of Ta2O5, BaCO3 and any dopant, being sure wiping spatula and inside of balance with acetone and kim-wipes between each chemical.
5. Place powders in milling bottle with MgO.
6. Add the grinding media labeled “BMT mix” and fill with acetone, ball mill for 24 hours.
7. Clean strainer in ultrasonic cleaner and rinse it with D.I. water and then acetone, then place it over a clean beaker.
8. Carefully pour powder slurry into a clean beaker and clean the grinding media over the strainer using acetone.
9. Be sure the media is clean before storing, rinsing with dilute HNO3 if necessary (not over the slurry beaker).
10. Dry sample on the hot plate using a magnetic stir bar to prevent settling (60°C & 200 RPM).
11. When all of the acetone is evaporated, place the beaker in drying oven for at least 24 hours.
12. Grind the powder with agate mortar and pestle as finely as possible.
13. Weigh platinum crucible empty, add powders and weigh again, record weights.
14. Calcine powders in an uncovered platinum crucible at 1300°C for 5 hours with a heating and cooling rate of 10°C/min.
15. Record the weight of the powder and crucible; calculate the weight loss.
16. Ball mill the powder with acetone in grinding media label “BMT Calcined” for 48 hours.
17. Remove ball media and dry sample as before.
18. When all acetone has evaporated, place beaker in drying oven for at least 24 hours.
19. Grind the powder with agate mortar and pestle as finely as possible.
**Uniaxial Pressing**

1. Clean ½” diameter die with isopropyl alcohol and kim-wipes.
2. Apply lubricant with swab to the contact surfaces of the die and top punch, but never on the face of the punch.
3. Weigh out 3.92 g of sample powder.
4. Place bottom punch in die, pour in powder, shake to settle it.
5. Slowly insert top punch.
6. Slide into press, being sure to hold bottom punch in place.
7. Close valve.
8. Increase pressure to 1,000 lbs, hold 5-10 secs, open valve slowly.
9. Remove bottom punch.
10. Invert die (with top punch in place) and set with foam under sides.
11. Push die down to force pellet up.
12. Remove pellet with tweezers.
13. Apply lubricant to die with new swab as needed (i.e. every other time).
14. Repeat steps 3-14 with rest of sample powder.
15. Clean die and punches with isopropyl alcohol.

**Cold Isostatic Pressing**

1. Open and rip rim off of non-lubricated condoms (# condoms = ½ # pellets).
2. Place pellet in bottom, remove all air with vacuum and knot.
3. Use 2 pellets for each condom; tie in separate compartments.
5. Put tube into oil in reactor (watch out for splash).
6. Settle back-up ring (oil should be ¼ inch below this level).
7. Open cover vent valve (allows excess oil to drain, use dish to catch).
8. Rotate cover by hand, holding center up with bar so it does not drop.
9. Tighten cover last few turns with bar, allowing time for oil to drain, and further tightening as this happens (don’t use too much force).
10. Once cover is snug, back it off 10-15 degrees to ease removal.
11. Close vent on cover when draining oil slows or stops.
12. Close vent valve on pump.
13. Open air inlet valve on pump.
14. Wait for gauge to read 40,000 psi.
15. Close inlet valve on pump.
16. Open vent valve on pump very slowly.
17. Open lid, remove tube and drain oil from tube.
18. Remove condoms from tube and squeeze off excess oil.
19. Over newspaper cut condoms open and remove pellets with tweezers, being careful to keep oil off pellets.

**Sintering**

1. Clean Pt crucible with acetone.
2. Weigh each pellet and record weight.
3. It is possible to sinter 4 pellets in a Pt crucible.
4. Note the location of each pellet using a reference point.
5. Cover pellets with smaller Pt crucible.
6. Sinter for specific amount of time and temperature.
7. Weigh and record weight of each pellet.
C.1 Network analyzer preparation

- A frequency sweep is a plot in the log magnitude format. This is indicated on the upper left of the network analyzer screen. To change the format to log mag, press FORMAT* on the panel then choose LOG MAG on the screen. There are two channels to observe the measurement. For convenience choose channel 1 for $S_{11}$ (for Q factor) and $S_{21}$ (for dielectric constant) measurements and channel 2 for the Smith chart. The channel shown on the screen can be changed using CH1 or CH2 buttons on the panel.
- Change the number of the data points to 1601: for maximum resolution.
  - Press MENU on the panel.
  - Choose NUMBER OF POINTS on the screen: When this is not 1601, change it to 1601 using the arrow keys on the panel.
- Press USER PRESET on the panel when something is strange or you don’t know what happened. Then start over.
- Refer “Operating manual for HP8722C Network Analyzer” for more detailed information about the network analyzer.
- Wear the wrist strap when you use the network analyzer for static free work.

* * Buttons on the network analyzer are highlighted.
C.2 Procedure for the Courtney method

The Courtney method is used for measuring the dielectric constant using the Courtney cavity (smaller brass cavity) and Courtney program. This is a transmission method and therefore uses 2 cables. Adaptors (one end is male and the other end is female (Wiltron) always attached to cavity) should be connected between the cavity and both cables. Thus the connection should be, from port 1, cable, gold connector (HP), adapter (wiltron), cavity, adapter (wiltron), gold connector (HP), and cable to port 2.

- Check channel and format of the screen. If they are not log magnitude format with channel 1 go to network analyzer preparation step (section C.1).
- Prepare network analyzer for the dielectric constant measurement.
  - Press MEAS on the panel.
  - Choose $S_{21}$ on the screen.
- Set scale of screen as 5dB.
  - Press SCALE REF on the panel.
  - Choose SCALE/DIV on the screen.
  - Hit 5 and x1 on the keypad.
- Set reference value as –15 dB.
  - Press SCALE REF on the panel.
  - Choose REFERENCE VALUE on the screen.
  - Enter -1 5 and x1 on the keypad.
  - Choose REFERENCE POSITION on the screen.
  - Move reference line using arrow keys on the panel.
  - Move small arrow on left side to the middle f the screen.
- Choose adequate frequency range (usually 1 to 10 GHz).
  - Press START on the panel and enter 1 then G/n on the keypad for setting a start frequency.
  - Press STOP on the panel and enter 1 0 on keypad then press G/n for setting a start frequency.
- Load the calibration from disk or manually calibrate the network analyzer. Before starting the calibration, make sure everything such as data points, span range, channel, and etc. is set up. You can recall it from the disk when you have a stored calibration for same frequency range (section C.6).
  - Disconnect cables from the cavity: Wiltron adaptor should be connected to cable.
  - Press CAL on the panel: Calibration menu will show up the screen.
    - Be sure that the CAL KIT is 2.92 mm: if it is not set up for 2.92 mm change it.
      - Choose CAL KIT on the screen.
      - Choose SELECT CAL KIT on the screen.
      - Choose 2.92 mm on the screen.
      - Choose RETURN on the screen: Calibration menu will show up the screen.
    - Choose CALIBRATION MENU on the screen.
    - Choose RESPONSE on the screen.
o Connect standard (Wiltron) (Both ends are male, kept in a plastic box in front of a network analyzer) to both cables with adaptors that were mentioned in C.2. The connections are now cable, from Port 1, gold connector (HP), adapter (wiltron), standard, adapter (wiltron), gold connector (HP), and cable to Port 2.
  o Choose **THRU** on screen: THRU screen will be underlined after a sweep.
  o Choose **DONE** on screen: Options for store calibration either on a disk or on a network analyzer will show up on the screen.

• Save calibration on a disk. (see Section C.4)
• Put a sample into the Courtney cavity.
  o Sample should be in center of the cavity.
  o Close lid and gently tighten three screws evenly.
  o Make sure that the lid is not tilted.
  o Screen now shows frequency spectrum. (Figure 42)
• Mark TE-011 mode peak.
  o Below 15 dB is not considered a peak. The first peak (above 15 dB) is the HE14 mode. Typical $S_{21}$ spectrum of the ceramic sample is shown in Figure C-1. It shows several resonant peaks. The TE-011 mode is the second lowest peak. Consider the one that is above the reference line (red line on the screen), -15dB, as a peak.
    ▪ Press **MKR** on the panel: Another menu will be shown.
    ▪ Choose **all OFF** on the screen: all markers will be erased.
    ▪ Choose **MODE MENU** on the screen: another menu will be shown.
    ▪ Choose **REF=1** on the screen.
    ▪ Point marker on the second lowest peak using the knob.
• Center TE-011 mode peak on the screen.
  o Press **MKR FCTN** on the panel: another menu will be shown.
  o Choose **MARKER MENU** on the screen.
  o Choose **MARKER CENTER** on the screen: Ignore the caution message.
• Reduce sweep span to 100 MHz.
  o Press **SPAN**: “span” will be shown on the left upper screen.
  o Press **1 0 0** and **M/μ** on keypad: Spectrum with 0.1 GHz span will be shown.
• Calibration (Repeat Cal step).
  o Calibrate a network analyzer again for this frequency span (0.1 GHz span), although a larger frequency span can be used if necessary.
• Save calibration.
• Connect cables to the cavity: Do not forget to connect Wiltron adaptor between cable and cavity. Connection should be cable, from Port 1., gold connector (HP), adapter (wiltron), cavity, adapter (wiltron), gold connector (HP), and cable to Port 2.
• Measure $S_{21}$ (insertion loss)
  o Press **MKR FCTN** on the panel: Another menu will be shown.
  o Choose **MAX** on the screen.
  o Choose **BANDWIDTH MENU** on the screen.
  o Choose **BW MEASURE** on the screen.
  o Choose **BANDWIDTH VALUE** on the screen.
  o Press -3 and **x1** on the keypad: This sets the BW value as –3dB.
• Read values for BW, loss ($S_{21}$), and cent (resonant frequency) from the screen, or see optional step following.
  o Press DISPLAY on the panel
  o Choose DATA $\rightarrow$ MEMORY on the screen.
  o Choose MEMORY on the screen: Now sweep stop.
  o Press MKR FCTN on the panel: Another menu will be shown.
  o Choose MAX on the screen.
  o Read values.
• (OPTIONAL) Print the values from screen.
  o Press COPY
  o Choose PRINT (66)
• Use Courtney program for calculating the dielectric constant.
• Change display mode to actual data.
  o Press DISPLAY on the panel.
  o Choose DISPLAY:DATA on the screen.
C.3 Courtney program

Use computer on the wooden desk (386 with windows 3.1).

- Change directory to C:/COURTNEY/COURTNEY/ on DOS mode.
- Run COURTNEY program.
  - Type COURTNEY then ENTER.
  - Hit ENTER again to continue the program: screen will be changed.
- Enter the values for program.
  - Enter values for pressure, temperature, humidity, and ambient dielectric constant as shown in figure 43. For diameter, height, resonance frequency, 3dB bandwidth, and insertion loss use the actual data for sample that you want to measure.
- Calculate the dielectric constant
  - Hit “Esc” key: Result will be shown as figure 44.
  - Hit “F9” key to return to previous display.
C.4 Save calibration

The calibration can be saved on a floppy disk or on the network analyzer. After finishing the calibration, options for save will be shown on a screen.

- Choose **STORE TO DISK** on the screen.
- Another menu will be shown.
- Initialize a disk: This is only for a new disk. **When you are using initialized disk skip this step.**
  - Insert a new disk into drive 0.
  - Choose **DEFINE, INIT, PURGE** on the screen.
  - Choose **INITIALIZE DISK** on the screen.
  - Choose **YES**: wait until the drive stops.
  - Choose **RETURN**: Another menu will be shown.
- Label file name
  - Choose **STORE TO DISK** on the screen.
  - Choose **TITLE FILE** on the screen.
  - Choose one of **TITLE**s on the screen.
  - Choose **ERASE TITLE** on the screen.
  - Point to the letter using the knob and press **SELECT LETTER** on the screen: Repeat until the name is finished.
  - Choose **DONE** on the screen: Menu will be changed.
  - Choose **RETURN** twice on the screen.
  - Choose **STORE** (what you named it): Wait until the drive stops. When disk does not have enough space to store, caution message will pop up. Then insert a new disk and initialize first, then store the calibration.
C.5 Procedure for Q factor measurement

This method is called single port reflection technique. Only one cable port (port 1) is required. It uses a silver plated aluminum cavity. For this method, only one cable and HP connector are used, i.e. no Wiltron adaptors were used with single cable. Thus connection should be, from Port 1, cable, yellow connector (HP), and cavity.

- Check channel and format of the screen. If they are not log magnitude format with channel 1 go to network analyzer preparation step (chapter C.1).
- Prepare network analyzer for the Q factor measurement.
  - Press MEAS on the panel.
  - Choose $S_{11}$ on the screen.
- Set scale of screen as 2dB.
  - Press SCALE REF on the panel.
  - Choose SCALE/DIV on the screen.
  - Enter 2 and x1 on the keypad.
- Set reference value as 0 dB.
  - Press SCALE REF on the panel.
  - Choose REFERENCE VALUE on the screen.
  - Enter 0 and x1 on the keypad.
  - Choose REFERENCE POSITION on the screen.
  - Move reference line using arrow key on the panel to (2) two lines from the top.
- Place a sample into a cavity.
  - It should be on the center of a styrofoam support.
  - Close lid gently.
- Choose adequate frequency range: It depends on the size of the ceramic sample and its dielectric constant. For the sample prepared in this lab, it starts from 1GHz.
  - Press START and enter 1 and G/n on the keypad. This changes the starting frequency to 1 GHz.
  - Set span as 100 MHz.
    - Press 0, 1, and G/n on the keypad: Spectrum with 0.1 GHz span will be shown.
- Sweep the sample with 100 MHz span from 1GHz until you find the peak.
  - Keep the same span.
  - Press CENTER and enter 1 and G/n on the keypad. This changes the center frequency to 1 GHz.
  - After one sweep and no peaks are observed, increase the center frequency by 100 MHz.
    - Press CENTER and enter new value using keypad then press G/n for a new center frequency.
  - Check TE-01$\delta$ mode.
    - Whenever you find the peak, open the lid of the cavity a little bit. If the peak disappears, keep increasing the center frequency by 100MHz until you find the peak that stays at the same position with the lid open. When a piece of metal is inserted through a hole into the cavity the peak should
shift slightly up in frequency. This confirms that this peak is TE-01δ mode.

- Mark the TE-01δ mode peak.
  - Press **MKR** on the panel: Another menu will be shown.
  - Choose all **OFF** on the screen: All markers will be erased.
  - Choose **MARKER=1** on the screen
  - Point marker to the resonant peak identified in last step.

- Center TE-01δ mode peak on the screen.
  - Press **MKR FCTN** on the panel: Another menu will be shown.
  - Choose **MARKER MENU** on the screen.
  - Choose **MARKER CENTER** on the screen: Ignore the caution message.

- Reduce span to 1 MHz.
  - Press **SPAN**: “span” will be shown on the left upper screen.
  - Press **1** and **M/μ** on keypad: Spectrum with 1 MHz span will show up (entire peak must be on the screen). If the range is too small to include the whole peak, keep increasing the span by 1 MHz increments until it fits.

- Calibration.
  - If you have a stored calibration for the frequency then recall it from the disk. (See section C.6)
  - If you don’t have a stored calibration you must calibrate. Disconnect the cable from the cavity: Yellow HP connector should be with a cable.
  - Press **CAL**.
    - Choose **CALIBRATION MENU** on screen.
    - Choose **S11 1-PORT** on the screen.
    - Connect the three standards (yellow HP, kept in a plastic box in front of a network analyzer) to cable the one by one. There are three standards (OPEN, CLOSE, and LOAD). They may be done any order.
      - Connect the CLOSE standard to the cable. (The close and short standards are labeled, the load is the smaller thinner standard that isn’t labeled.)
      - Choose **CLOSE** on the screen: **CLOSE** will be underlined after a sweep.
    - Connect the OPEN standard to the cable.
    - Choose **OPEN** on the screen: **OPEN** will be underlined after a sweep.
    - Connect the LOAD standard to cable.
    - Choose **LOAD** on the screen: a different menu will show up.
    - Choose **BROADBAND** on the screen: Broadband will be underlined after a sweep.
    - Choose **DONE** on the screen for returning to Cal menu.
      - Choose **DONE 1-PORT CAL** on the screen: the analyzer will calculate the calibration. The save menu will show up.
  - Save calibration on a disk. (section C.4)
  - Adjust the coupling probe inside of the cavity.
Press CH2: Smith chart (channel 2) will show up instead of the log magnitude format. The coupling loop will show up on the Smith chart. The coupling loop should be touching the big circle of the Smith chart at a point where resistance is infinite (Figure 44), i.e. the edge.

If a loop is out of the center of the Smith chart move it into center.
- Press SCALE REF.
- Choose PHASE OFFSET on the screen.
- Move the coupling loop using the knob to center circle.
- Move the probe a little bit with extra caution.
- Close the lid after adjusting the probe.
- Observe the coupling loop on the screen of the network analyzer.
- Adjust the coupling probe until the coupling loop of on the Smith chart overlaps the 50 ohms circle (Figure C-4).

Display both channels simultaneous
- Press DISPLAY
- Choose DUAL CHAN on the screen: Log magnitude format (channel 1) will be above the Smith chart (channel 2). (Figure 45)

Measure $S_{11}$ (return loss)
- Press CH1 to control the log magnitude format.
- Press MKR FCTN on the panel: Another menu will be shown.
- Choose MIN on the screen as soon as one sweep is done. From this step to choosing “data→memory” step should be done within one full sweep. Otherwise minimum point on channel 1 may be different value from channel 2, since the return losses are slightly different from each sweep.
- At the same time read the resistance on the channel 2: It is highlighted in the Figure A-5.
- Press DISPLAY
- Choose DATA→MEMORY on the screen.
- Choose MEMORY on the screen: Now sweep stops.
- Determine under/over coupling.
  - If the resistance at the minimum point of the resonant peak is less than 50 ohms, it is an overcoupling.
  - If the resistance at the minimum point of the resonant peak is larger than 50 ohms, it is an undercoupling.

Measure the unloaded Q factor from the network analyzer.
- Press MKR FCTN on the panel.
- Choose BANDWIDTH MENU on the screen.
- Choose BW MEASURE on the screen.
- Read return loss (RL) from the screen.

Calculate modified return loss (RL’).
- Use a computer sit on right hand side (Pentium with Windows 98).
- Double click “Couple” from the desktop menu: Icon is on the screen.
- Enter value of RL into spreadsheet (Figure 46).
- Hit ENTER.
- Modified return loss (RL’) will be calculated.

Measure the unloaded Q factor from the network analyzer.
114

- Press **MKR FCTN** on the panel.
- Choose **BANDWIDTH MENU** on the screen.
- Choose **BANDWIDTH VALUE** on the screen.
- Enter the RL' value using the keypad then press **x1**.
- Read Q value.
- Print out the result.
  - Press **COPY** on the panel.
  - Choose **PRINT** on the screen.
- Reset for next measurement
  - Display log magnitude format on the screen.
    - Press **CH1** on the panel.
    - Press **DISPLAY** on the panel.
    - Choose **DUAL CHAN** on the screen (off).
  - Change display mode to actual data.
    - Press **DISPLAY** on the panel
    - Choose **DISPLAY:DATA** on the screen.
A.6 Recall the calibration file from a disk.

- Insert a disk into the drive 0.
- Press **RECALL**: Menu will show up the screen.
- Choose **LOAD FROM DISK** on the screen: Another menu will show up on the screen.
- Choose **LOAD FILE TITLES** on the screen: It will show 5 titles at a time from the disk. Repeat until you see the filename that you want to load on the network analyzer.
- Choose filename that you want on the screen.
**Figure 47** The $S_{21}$ spectrum ranging from 1 GHz to 10 GHz.
Figure 48 The screen on the Courtney program where the actual data is entered so the program can calculate the dielectric constant.
The test specimen identification is Defaults.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Specimen diameter (cm)</td>
<td>1.000</td>
</tr>
<tr>
<td>Specimen height (cm)</td>
<td>0.500</td>
</tr>
<tr>
<td>Coeff. thermal exp. (ppm)</td>
<td>0</td>
</tr>
<tr>
<td>Ambient temperature (°C)</td>
<td>25.0</td>
</tr>
<tr>
<td>Ambient dielectric const.</td>
<td>1.00053</td>
</tr>
</tbody>
</table>

The input data and results for the TE-011 mode are:

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Resonance frequency (GHz)</td>
<td>7.0000</td>
</tr>
<tr>
<td>3dB bandwidth (GHz)</td>
<td>0.00400</td>
</tr>
<tr>
<td>Insertion loss (dB)</td>
<td>25.0</td>
</tr>
<tr>
<td>Loss tangent (ppm)</td>
<td>400</td>
</tr>
<tr>
<td>Dielectric constant</td>
<td>35.65</td>
</tr>
</tbody>
</table>

The maximum frequency for this mode is 29.97 GHz.

**Figure 49** An example of the results obtained from the Courtney Software.
Figure 50 The resonant coupling circle on the Smith chart for measuring the Q-factor using the single-port reflection technique
Figure 51  An example of the results obtained from the single-port reflection technique for measuring the Q-factor.
BIBLIOGRAPHY
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